

Technische Universität Graz

Optical Properties of Metals

AC Conductivity

For constant voltage, conductors conduct and insulators don't.

For low ac voltages in a conductor, electric field and the electron velocity are inphase, electric field and electron position are out-of-phase.

For low ac voltages in an insulator, electric field and the electron position are inphase, electric field and electron velocity are out-of-phase.

At high (optical) frequencies the in-phase and out-of-phase component of the response is described by the dielectric function.

Conductivity / Dielectric function

Harmonic dependence
$$v = v(\omega)e^{i\omega t}$$
, $x = x(\omega)e^{i\omega t}$, $E = E(\omega)e^{i\omega t}$

$$v(\omega) = i\omega x(\omega)$$

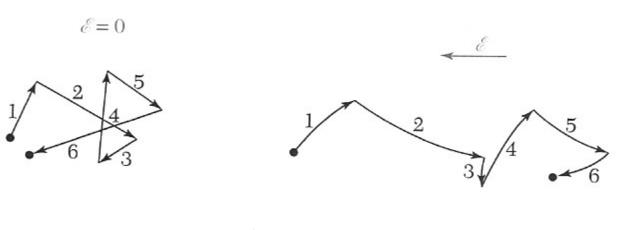
$$\chi(\omega) = \frac{P(\omega)}{\varepsilon_0 E(\omega)} = \frac{-nex(\omega)}{\varepsilon_0 E(\omega)} \qquad \sigma(\omega) = \frac{j(\omega)}{E(\omega)} = \frac{-nev(\omega)}{E(\omega)} = \frac{-i\omega nex(\omega)}{E(\omega)}$$

$$\chi(\omega) = \frac{\sigma(\omega)}{i\omega\varepsilon_0}$$

$$\varepsilon(\omega) = 1 + \chi = 1 + \frac{\sigma(\omega)}{i\omega\varepsilon_0}$$

Below about 100 GHz the frequency dependent conductivity is normally used. Above about 100 GHz the dielectric function is used (optical experiments).

Diffusive transport (low frequencies)



$$\begin{split} \vec{F} &= m\vec{a} = -e\vec{E} = m\frac{\vec{v}_d}{\tau_{\rm sc}} \\ &-\frac{e\tau_{\rm sc}}{m}\vec{E} = \vec{v}_d \\ &-\mu_e\vec{E} = \vec{v}_d \end{split} \qquad \qquad \sigma = ne\mu = \frac{ne^2\tau}{m} \end{split}$$

Diffusive metal

The current is related to the electric field

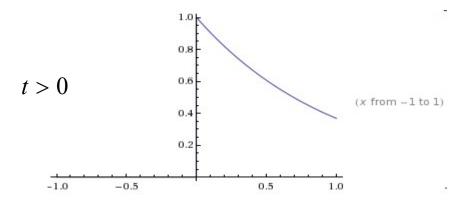
$$j_n = \sigma_{nm} E_m$$
 $v_n = -\mu_{nm} E_m$ Steady state solution

The differential equation that describes how the velocity changes in time is:

$$m\frac{dv(t)}{dt} + \frac{ev(t)}{\mu} = -eE(t)$$
Inertial term

The impulse response function:

$$g(t) = \frac{1}{m} \exp\left(\frac{-et}{\mu m}\right)$$



Diffusive metal

The differential equation is:

$$m\frac{dv(t)}{dt} + \frac{ev(t)}{\mu} = -eE(t)$$

Assume a harmonic solution $E(\omega)e^{i\omega t}$, $v(\omega)e^{i\omega t}$

$$\left(-\frac{i\omega m}{e} - \frac{1}{\mu}\right) v(\omega) = E(\omega)$$

$$\frac{v(\omega)}{E(\omega)} = \left(-\frac{i\omega m}{e} - \frac{1}{\mu}\right)^{-1} = -\mu \left(1 + i\omega\tau\right)^{-1} = \frac{-\mu \left(1 - i\omega\tau\right)}{1 + \omega^2\tau^2}$$

$$\sigma(\omega) = \frac{j(\omega)}{E(\omega)} = -ne\frac{v(\omega)}{E(\omega)} = ne\mu \left(\frac{1 - i\omega\tau}{1 + \omega^2\tau^2}\right) \qquad \tau = \frac{\mu m}{e}$$
 Scattering time

$$\sigma(\text{low }\omega) = ne\mu$$

$$\tau = \frac{\mu m}{e}$$
 Scattering time

$$\sigma(\text{high }\omega) = \frac{-ine^2}{\omega m}$$

Diffusive metal

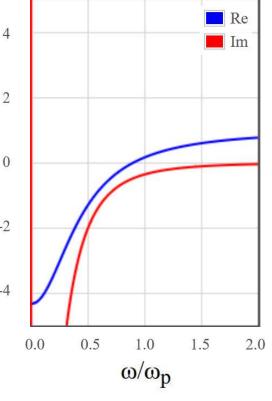
$$\chi(\omega) = \frac{\sigma(\omega)}{i\omega\varepsilon_0} = \frac{ne\mu}{i\omega\varepsilon_0} \left(\frac{1 - i\omega\tau}{1 + \omega^2\tau^2} \right)$$

$$\varepsilon(\omega) = 1 + \chi = 1 - \frac{ne\mu}{\omega\varepsilon_0} \left(\frac{\omega\tau + i}{1 + \omega^2\tau^2} \right)$$

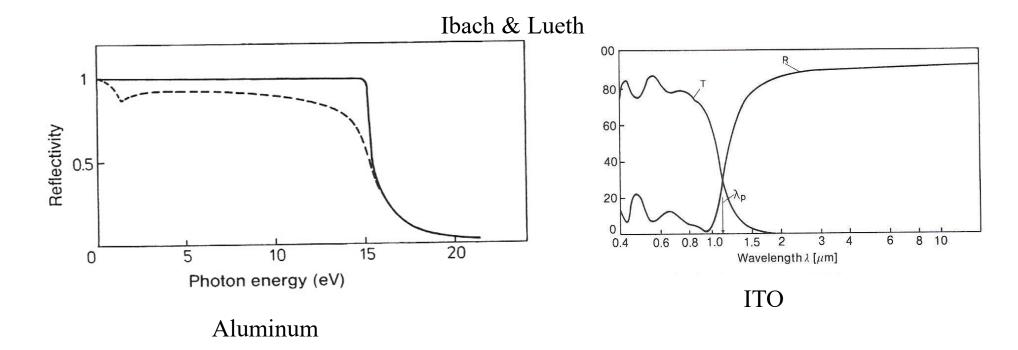
$$\varepsilon(\omega) = 1 - \omega_p^2 \left(\frac{\omega \tau^2 + i\tau}{\omega + \omega^3 \tau^2} \right) \qquad \omega_p^2 = \frac{ne^2}{m\varepsilon_0} \qquad \epsilon_r(\omega)^{-0}$$

Take the limit as τ goes to infinity

$$\varepsilon'(\omega) = 1 - \frac{\omega_p^2}{\omega^2}$$



low frequency metal / high frequency insulator

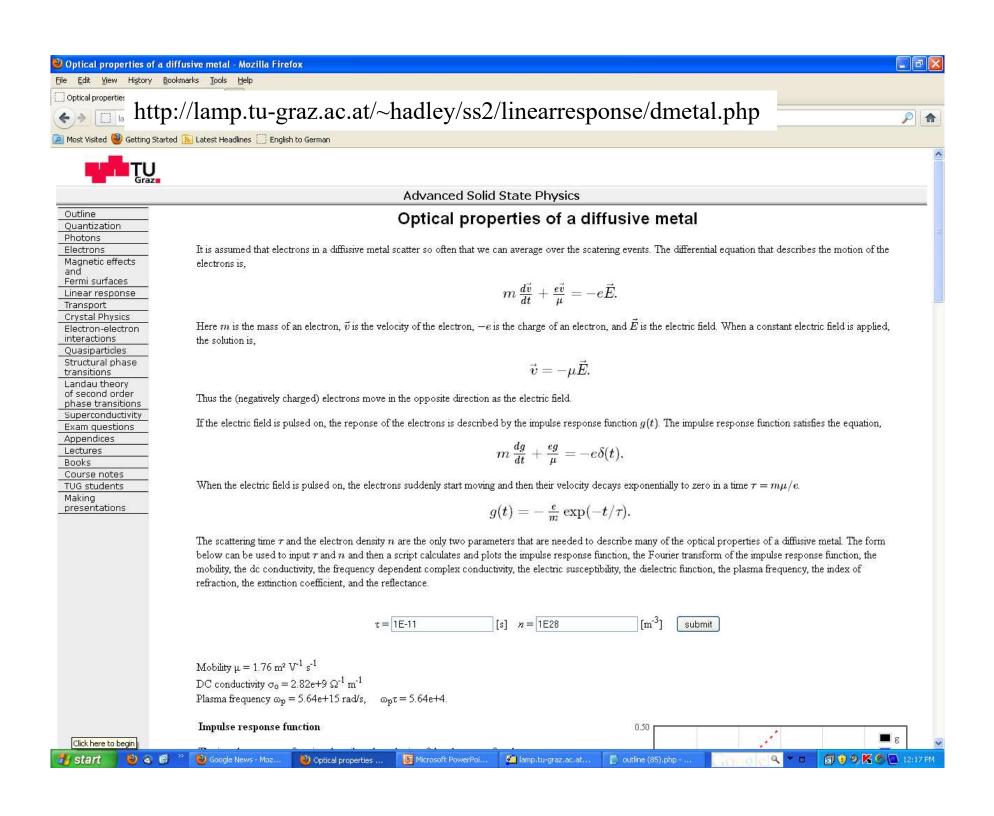


Conducting transparent contacts for LEDs and Solar cells

Windows that reflect infrared

Reflection of radio waves from ionosphere

$$\omega_p^2 \approx \frac{ne^2}{\varepsilon_0 m}$$





Advanced Solid State Physics

Outline

Quantization Photons

Electrons

Magnetic effects

Fermi surfaces

Linear response

Transport

Crystal Physics

Electron-electron interactions

Quasiparticles Structural phase

transitions

Landau theory of second order phase transitions

Superconductivity
Exam questions

Appendices

Lectures

Books

Course notes

TUG students

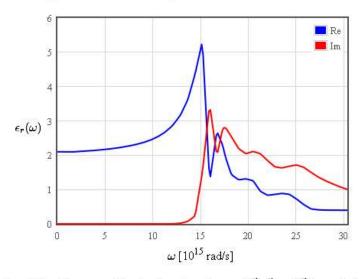
Making presentations Cu Si SiO2 diamond

The optical properties of SiO₂ (glass)

nanophotonics.csic.es

Dielectric function

The relative dielectric constant describes the relationship between the electric displacement \vec{D} and the electric field \vec{E} , $\vec{D} = \epsilon_r \epsilon_0 \vec{E} = \vec{P} + \epsilon_0 \vec{E}$.



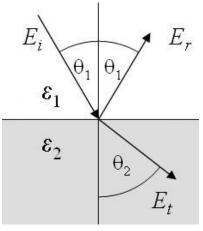
There are two conventions for dielectric function. Either it is assumed that the time dependence of \vec{D} , \vec{P} , and \vec{E} is $\exp(-i\omega t)$ and the plot of the dielectric function looks as it is shown above, or it is assumed that the time dependence of \vec{D} , \vec{P} , and \vec{E} is $\exp(i\omega t)$ and the imaginary part of the has the opposite sign as in the plot above. Here we will assume a time dependence of $\exp(-i\omega t)$.

Electric susceptibility

The electric susceptibility χ_E describes the relationship between the polarization \vec{P} and the electric field \vec{E} , $\vec{P} = \epsilon_0 \chi_E \vec{E}$.

$$\chi_{E}=\epsilon_{r}-1$$

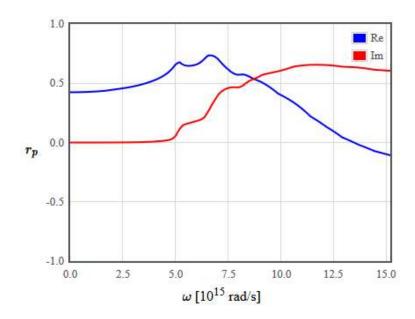
Ellipsometry

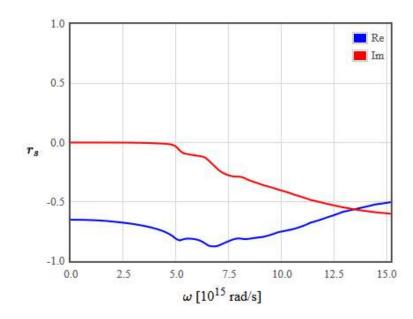


$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

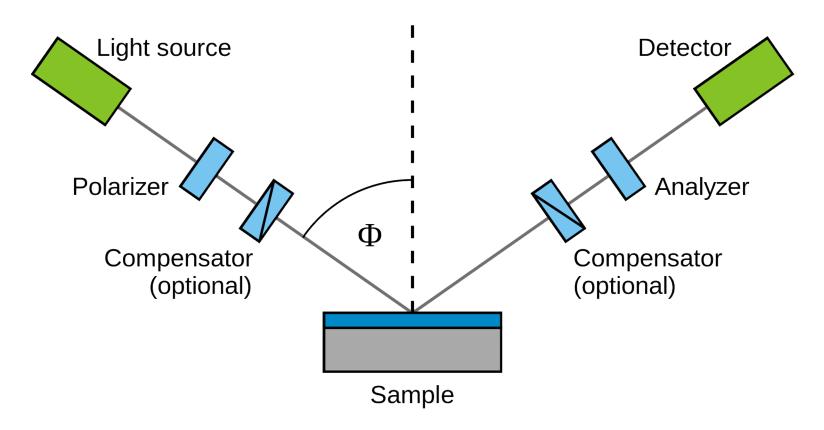
$$r_p = \frac{E_{rp}}{E_{ip}} = \frac{\sqrt{\epsilon_2} \cos \theta_1 - \sqrt{\epsilon_1} \cos \theta_2}{\sqrt{\epsilon_2} \cos \theta_1 + \sqrt{\epsilon_1} \cos \theta_2}$$

$$r_s = rac{E_{sr}}{E_{si}} = rac{\sqrt{\epsilon_2}\cos heta_2 - \sqrt{\epsilon_1}\cos heta_1}{\sqrt{\epsilon_1}\cos heta_1 + \sqrt{\epsilon_2}\cos heta_2}$$





Ellipsometry

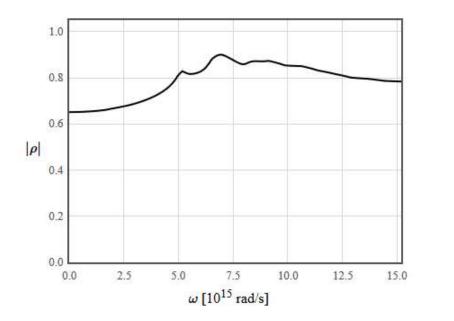


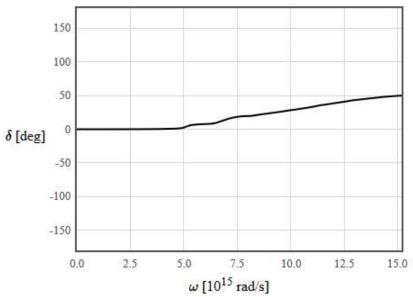
Ellipsometry measures the change of polarization upon reflection. The measured signal depends on the thickness and the dielectric constant.

http://en.wikipedia.org/wiki/Ellipsometry

Ellipsometry

$$ho=rac{r_p}{r_s}=|
ho|e^{i\delta}$$





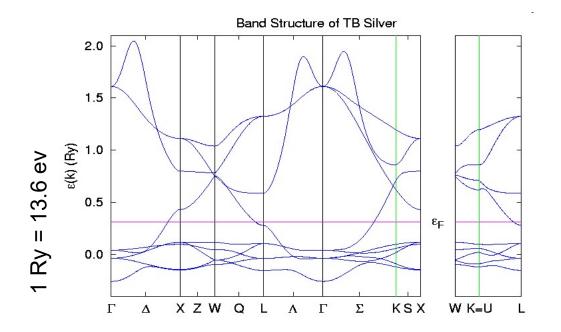
The ratio of the two reflected polarizations is insensitive to instabilities of light source or atmospheric absorption.

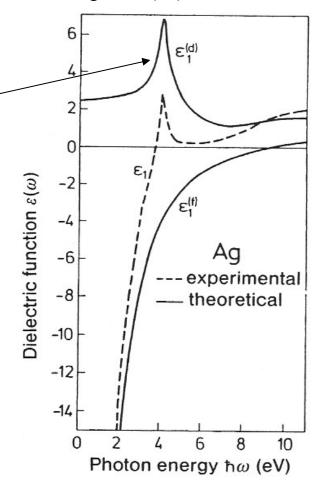
Intraband transitions

When the bands are parallel, there is a peak in the absorption (ϵ ")

 $\hbar\omega = E_c(\vec{k}) - E_v(\vec{k})$

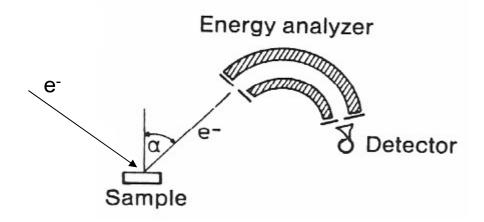
Intraband (d-band) absorption



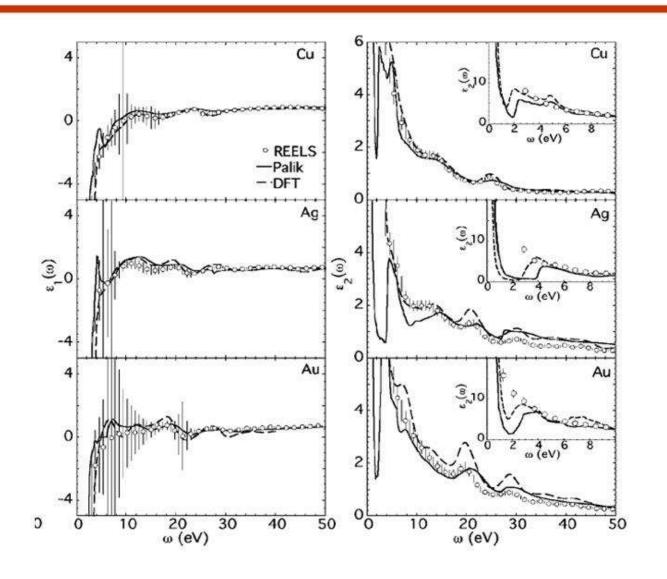


Ibach & Lueth

Reflection Electron Energy Loss Spectroscopy



Fast electrons moving through the solid generate and a time dependent electric field. If the polarization moves out of phase with this field, energy will be lost. This is detected in the reflected electrons.



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