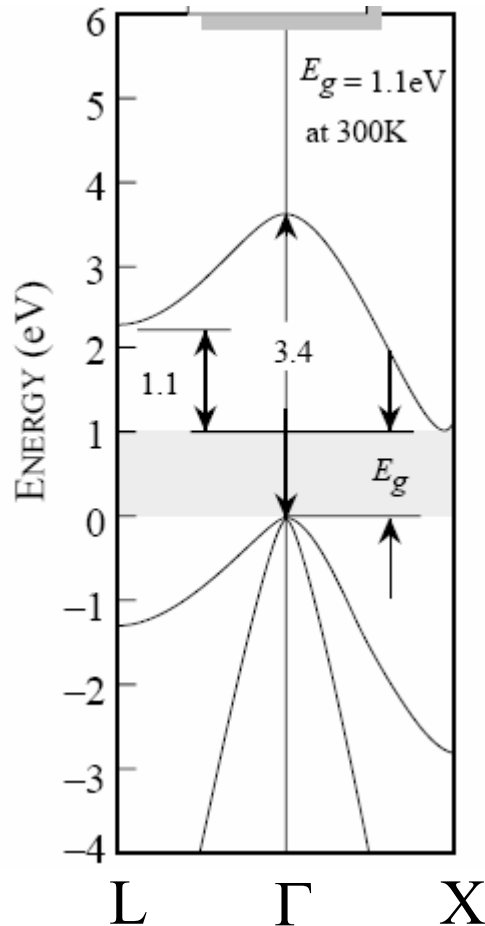


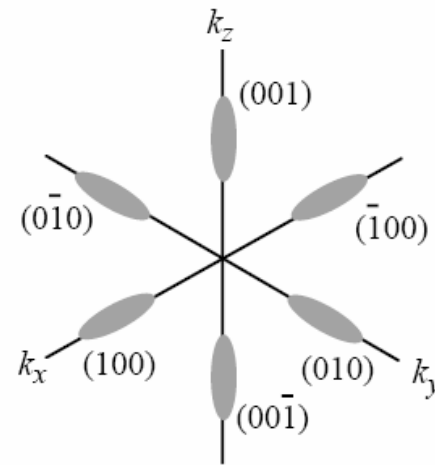
# Intrinsic Semiconductors

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# Density of electrons in the conduction band



$$n = N_c(T) \exp\left(\frac{\mu - E_c}{k_B T}\right) = \frac{2D_v}{\sqrt{\pi}} (k_B T)^{3/2} \exp\left(\frac{\mu - E_c}{k_B T}\right)$$

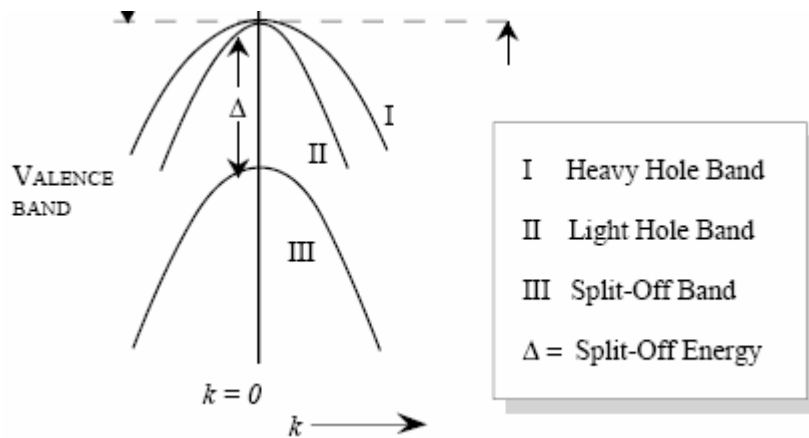


$$n = N_c \left(\frac{T}{300}\right)^{3/2} \exp\left(\frac{E_F - E_c}{k_B T}\right)$$

Properties	Si	Ge	GaAs
Bandgap $E_g$	1.12 eV	0.66 eV	1.424 eV
Effective density of states in conduction band (300 K) $N_c$	$2.78 \times 10^{25} \text{ m}^{-3}$	$1.04 \times 10^{25} \text{ m}^{-3}$	$4.45 \times 10^{23} \text{ m}^{-3}$

# Density of electrons in the conduction band

## Density of holes in the valence band



$$n = N_c \left( \frac{T}{300} \right)^{3/2} \exp \left( \frac{\mu - E_c}{k_B T} \right)$$

$$p = N_v \left( \frac{T}{300} \right)^{3/2} \exp \left( \frac{E_v - \mu}{k_B T} \right)$$

Properties	Si	Ge	GaAs
Bandgap $E_g$	1.12 eV	0.66 eV	1.424 eV
Effective density of states in conduction band (300 K) $N_c$	$2.78 \times 10^{25} \text{ m}^{-3}$	$1.04 \times 10^{25} \text{ m}^{-3}$	$4.45 \times 10^{23} \text{ m}^{-3}$
Effective density of states in valence band (300 K) $N_v$	$9.84 \times 10^{24} \text{ m}^{-3}$	$6.0 \times 10^{24} \text{ m}^{-3}$	$7.72 \times 10^{24} \text{ m}^{-3}$
Effective mass electrons $m^*/m_0$	$m_l^* = 0.98$ $m_t^* = 0.19$	$m_l^* = 1.64$ $m_t^* = 0.082$	$m^* = 0.067$
Effective mass holes $m^*/m_0$	$m_{lh}^* = 0.16$ $m_{hh}^* = 0.49$	$m_{lh}^* = 0.044$ $m_{hh}^* = 0.28$	$m_{lh}^* = 0.082$ $m_{hh}^* = 0.45$
Crystal structure	diamond	diamond	zincblende
Density	2.328 g/cm <sup>3</sup>	5.3267 g/cm <sup>3</sup>	5.32 g/cm <sup>3</sup>
Atoms/m <sup>3</sup>	$5.0 \times 10^{28}$	$4.42 \times 10^{28}$	$4.42 \times 10^{28}$

# The electrical contribution to the thermodynamic properties of insulators depend on band edges

## Boltzmann approximation

The table below gives the contribution of electrons in intrinsic semiconductors and insulators to some thermodynamic quantities. These results were calculated in the Boltzmann approximation where it is assumed that the chemical potential lies in the band gap more than  $3k_B T$  from the band edge. The electronic contribution to the thermodynamic quantities are usually much smaller than the contribution of the phonons and thus the electronic components are often simply ignored.

	1-d	2-d
<b>Density of states</b> $m_e^*$ and $m_h^*$ are 'density of states' effective masses	$D(E) = \begin{cases} \frac{1}{\hbar\pi} \sqrt{\frac{2m_h^*}{(E_v - E)}} & E < E_v \\ 0 & E_v < E < E_c \\ \frac{1}{\hbar\pi} \sqrt{\frac{2m_e^*}{(E - E_c)}} & E > E_c \end{cases} \quad \text{J}^{-1} \text{m}^{-1}$	$D(E) = \begin{cases} \frac{m_h^*}{\hbar^2\pi} H(E_v - E) & E < E_v \\ 0 & E_v < E < E_c \\ \frac{m_e^*}{\hbar^2\pi} H(E - E_c) & E > E_c \end{cases} \quad \text{J}^{-1} \text{m}^{-2}$ $H(x) = 0 \text{ for } x < 0 \text{ and } H(x) = 1 \text{ for } x > 0$
<b>Density of states</b> $N_v$ and $N_c$ are the effective densities of states	$D(E) = \begin{cases} N_v(300) \sqrt{\frac{2}{300\pi k_B (E_v - E)}} & E < E_v \\ 0 & E_v < E < E_c \\ N_c(300) \sqrt{\frac{2}{300\pi k_B (E - E_c)}} & E > E_c \end{cases} \quad \text{J}^{-1} \text{m}^{-1}$	$D(E) = \begin{cases} \frac{N_v(300)}{300k_B} H(E_v - E) & E < E_v \\ 0 & E_v < E < E_c \\ \frac{N_c(300)}{300k_B} H(E - E_c) & E > E_c \end{cases} \quad \text{J}^{-1} \text{m}^{-2}$
<b>Density of electrons in the conduction band</b> $n = \int_{E_c}^{\infty} D(E) f(E) dE$	$n = \sqrt{\frac{m_e^* k_B T}{\hbar^2 \pi}} \exp\left(\frac{\mu - E_c}{k_B T}\right) \quad \text{m}^{-1}$ $= N_c \exp\left(\frac{\mu - E_c}{k_B T}\right)$	$n = \frac{m_e^* k_B T}{\hbar^2 \pi} \exp\left(\frac{\mu - E_c}{k_B T}\right) \quad \text{m}^{-2}$ $= N_c \exp\left(\frac{\mu - E_c}{k_B T}\right)$
<b>Density of holes in the valence band</b> $p = \int_{-\infty}^{E_v} D(E) (1 - f(E)) dE$	$p = \sqrt{\frac{m_h^* k_B T}{\hbar^2 \pi}} \exp\left(\frac{E_v - \mu}{k_B T}\right) \quad \text{m}^{-1}$ $= N_v \exp\left(\frac{\mu - E_c}{k_B T}\right)$	$p = \frac{m_h^* k_B T}{\hbar^2 \pi} \exp\left(\frac{E_v - \mu}{k_B T}\right) \quad \text{m}^{-2}$ $= N_v \exp\left(\frac{\mu - E_c}{k_B T}\right)$

## Intrinsic semiconductors

In the Boltzmann approximation, the density of states of a semiconductor is,

$$D(E) = \begin{cases} \frac{(2m_h^*)^{3/2}}{2\pi^2\hbar^3} \sqrt{E_v - E}, & \text{if } E < E_v \\ 0, & \text{if } E_v < E < E_c \\ \frac{(2m_e^*)^{3/2}}{2\pi^2\hbar^3} \sqrt{E - E_c}, & \text{if } E_c < E \end{cases}$$

Here  $m_e^*$  and  $m_h^*$  are the 'density of states effective masses' for electrons and holes. Usually in the literature, effective density of states at 300 K is given instead of the 'density of states effective masses'. The relationship between the two is,

$$m_h^* = \frac{\pi\hbar^2}{300k_B} \left( \sqrt{2}N_v(300) \right)^{2/3}$$

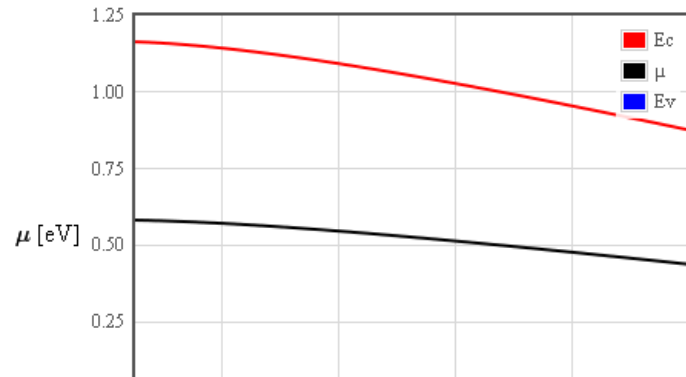
$$m_e^* = \frac{\pi\hbar^2}{300k_B} \left( \sqrt{2}N_c(300) \right)^{2/3}$$

In an intrinsic semiconductor, the density of electrons equals the density of holes,  $n = p = n_i = \sqrt{N_c \left( \frac{T}{300} \right)^{3/2} N_v \left( \frac{T}{300} \right)^{3/2} \exp\left( \frac{-E_g}{2k_B T} \right)}$ .

By setting the concentration of electrons equal to the concentration of holes, it is possible to solve for the chemical potential. The bandgap of most semiconductors is temperature dependent. The form below lets you input the temperature dependence of the bandgap. The bandgaps for some semiconductors can be loaded into the form with the buttons on the right.

$$n = N_c(300) \left( \frac{T}{300} \right)^{3/2} \exp\left( \frac{\mu - E_c}{k_B T} \right) = p = N_v(300) \left( \frac{T}{300} \right)^{3/2} \exp\left( \frac{E_v - \mu}{k_B T} \right)$$

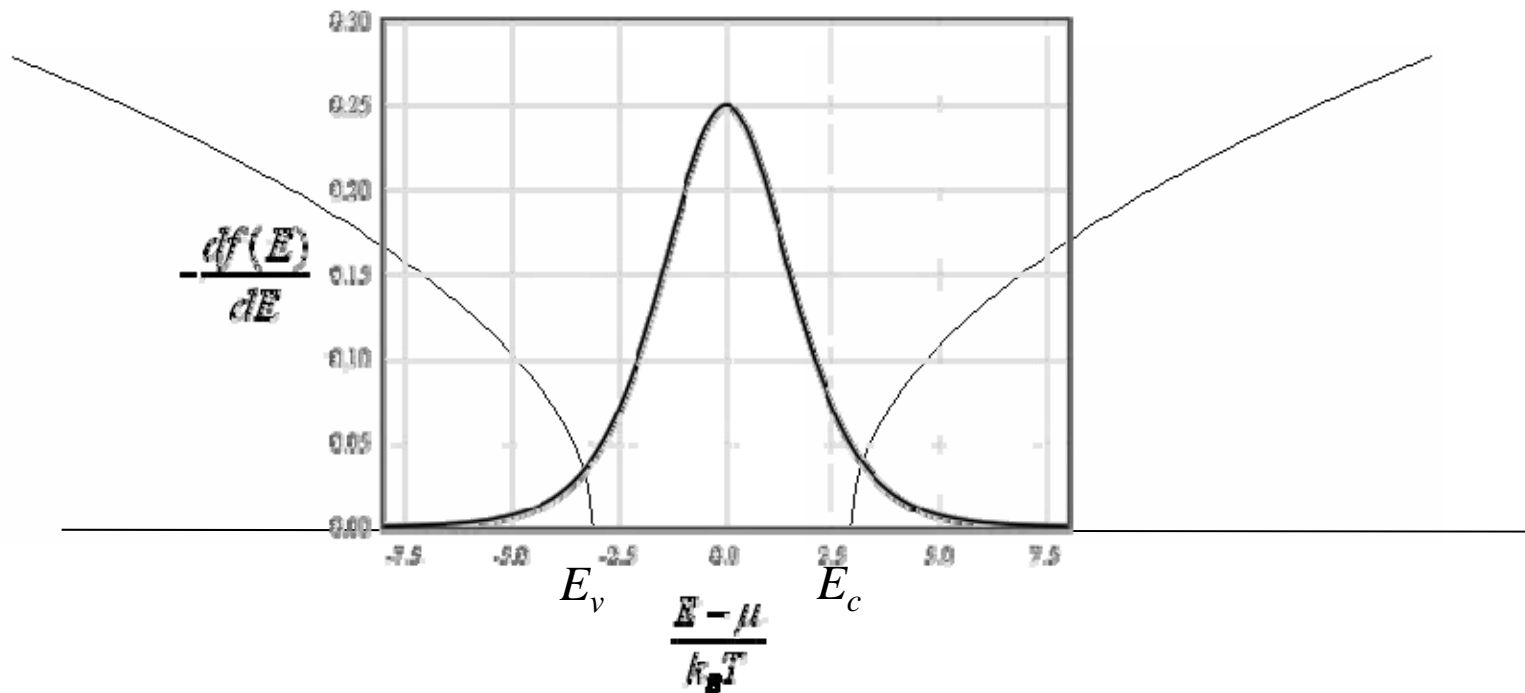
$$\mu = \frac{E_v + E_c}{2} + k_B T \ln\left( \frac{N_v(300)}{N_c(300)} \right)$$



$N_c(300 \text{ K}) =$	<input type="text" value="2.78E19"/>	1/cm <sup>3</sup>	Semiconductor <input type="button" value="Si"/> <input type="button" value="Ge"/> <input type="button" value="GaAs"/>
$N_v(300 \text{ K}) =$	<input type="text" value="9.84E18"/>	1/cm <sup>3</sup>	
$E_g =$	<input type="text" value="1.166-4.73E-4*T*(T+636)"/>	eV	
$T_1 =$	<input type="text" value="50"/>	K	
$T_2 =$	<input type="text" value="1000"/>	K	
<input type="button" value="Replot"/>			

# Narrow bandgap semiconductors

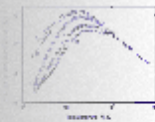
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Use the programs for metals for small bandgap semiconductors.

# Semiconductors

on NSM



Semiconductors

n, k database

InGaAsP

Equivalents

<a href="#">Si</a>	- Silicon	<a href="#">Ge</a>	- Germanium
<a href="#">GaP</a>	- Gallium Phosphide	<a href="#">GaAs</a>	- Gallium Arsenide
<a href="#">InAs</a>	- Indium Arsenide	<a href="#">C</a>	- Diamond
<a href="#">GaSb</a>	- Gallium Antimonide	<a href="#">InSb</a>	- Indium Antimonide
<a href="#">InP</a>	- Indium Phosphide	<a href="#">GaAs<sub>1-x</sub>Sb<sub>x</sub></a>	- Gallium Arsenide Antimonide
<a href="#">Al<sub>x</sub>Ga<sub>1-x</sub>As</a>	- Aluminium Gallium Arsenide		
<a href="#">AlN</a>	- Aluminium Nitride	<a href="#">InN</a>	- Indium Nitride
<a href="#">BN</a>	- Boron Nitride	<a href="#">GaN</a>	- Gallium Nitride

## We are going to add new data for:

<a href="#">Ga<sub>x</sub>In<sub>1-x</sub>As<sub>y</sub>Sb<sub>1-y</sub></a>	- Gallium Indium Arsenide Antimonide	<a href="#">Ga<sub>x</sub>In<sub>1-x</sub>P</a>	- Gallium Indium Phosphide
<a href="#">Ga<sub>x</sub>In<sub>1-x</sub>As</a>	- Gallium Indium Arsenide	<a href="#">Ga<sub>x</sub>In<sub>1-x</sub>Sb</a>	- Gallium Indium Antimonide
<a href="#">InAs<sub>1-x</sub>Sb<sub>x</sub></a>	- Indium Arsenide Antimonide	<a href="#">Ga<sub>x</sub>In<sub>1-x</sub>As<sub>y</sub>P<sub>1-y</sub></a>	- Gallium Indium Arsenide Phosphide
<a href="#">Si<sub>1-x</sub>Ge<sub>x</sub></a>	- Silicon Germanium	<a href="#">SiC</a>	- Silicon Carbide

## Effective Masses

### Electrons:

The surfaces of equal energy are ellipsoids.

$$m_l = 0.98m_0$$

$$m_t = 0.19m_0$$

Effective mass of density of states  $m_c = 0.36m_0$

There are 6 equivalent valleys in the conduction band.

$$m_{cc} = 0.26m_0$$

### Holes:

Heavy  $m_h = 0.49m_0$

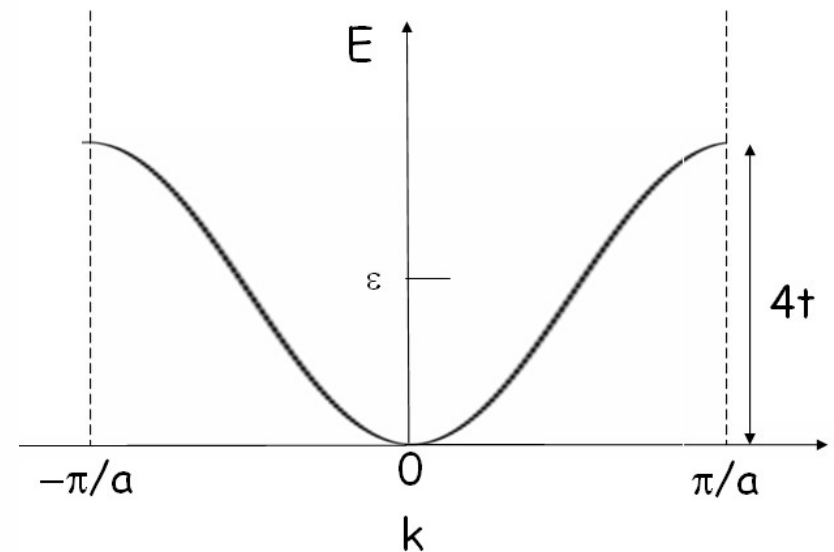
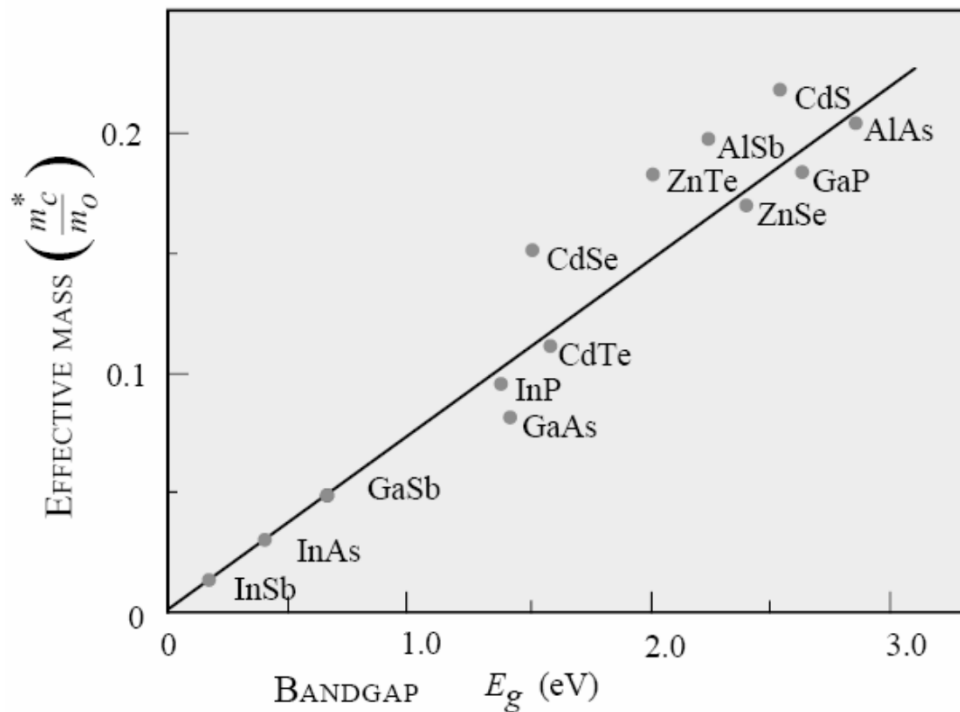
Light  $m_{lp} = 0.16m_0$

Split-off band  $m_{so} = 0.24m_0$

Effective mass of density of states  $m_v = 0.81m_0$



# Large gap -> large effective mass



$$E_k = \epsilon - 2t \cos(ka)$$

$$m^* = \frac{\hbar^2}{\frac{d^2 E}{dk^2}}$$

$$\frac{d^2 E}{dk^2} = 2ta^2$$

$$m^* \sim \frac{1}{t}$$

narrow bands -> large effective mass

# Measuring the effective mass

---

Cyclotron resonance  $\omega_c = \frac{eB}{m^*}$

Resonant absorption occurs when rf waves with the cyclotron resonance frequency are applied. This can be used to experimentally determine the effective mass.

Knowing the effective mass, the scattering time can be calculated from the measured conductivity.

$$\sigma = \frac{ne^2\tau_{sc}}{m^*}$$

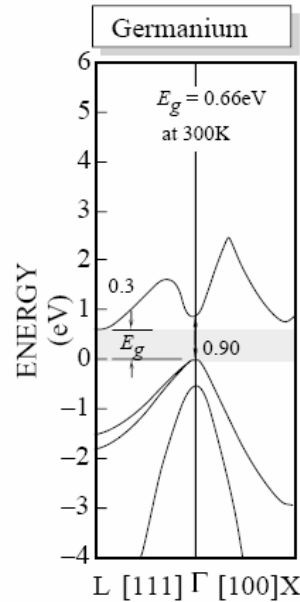
# Direct and indirect band gaps

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indirect bandgap

$$\Delta k \neq 0$$

phonons are emitted

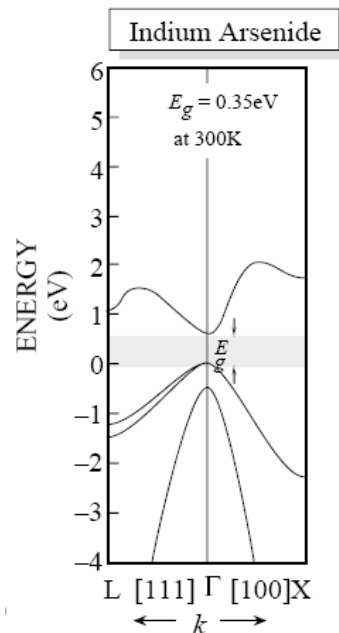


Momentum must be conserved when photons are absorbed or emitted.

direct bandgap:

$$\Delta k = 0$$

photons can be emitted



**TABLE 1 Common III-V materials used to produce LEDs and their emission wavelengths.**

Material	Wavelength (nm)
InAsSbP/InAs	4200
InAs	3800
GaInAsP/GaSb	2000
GaSb	1800
$Ga_xIn_{1-x}As_{1-y}P_y$	1100-1600
$Ga_{0.47}In_{0.53}As$	1550
$Ga_{0.27}In_{0.73}As_{0.63}P_{0.37}$	1300
GaAs:Er, InP:Er	1540
Si:C	1300
GaAs:Yb, InP:Yb	1000
$Al_xGa_{1-x}As:Si$	650-940
GaAs:Si	940
$Al_{0.11}Ga_{0.89}As:Si$	830
$Al_{0.4}Ga_{0.6}As:Si$	650
$GaAs_{0.6}P_{0.4}$	660
$GaAs_{0.4}P_{0.6}$	620
$GaAs_{0.15}P_{0.85}$	590
$(Al_xGa_{1-x})_{0.5}In_{0.5}P$	655
GaP	690
GaP:N	550-570
$Ga_xIn_{1-x}N$	340,430,590
SiC	400-460
BN	260,310,490

# Light emitting diodes

