15. Magnetism

Nov 21, 2019

Ferrimagnets

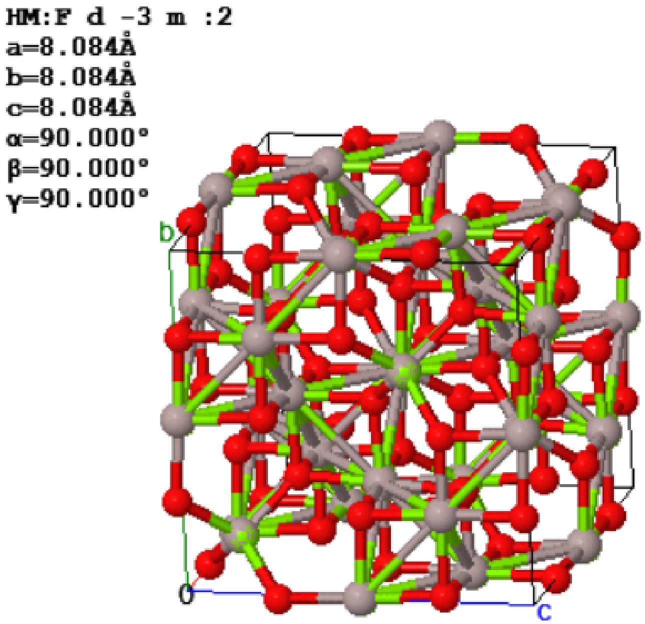
Magnetite Fe_3O_4
(Magneteisen)

Ferrites $MO{\cdot}Fe_2O_3$

M = Fe, Zn, Cd, Ni, Cu, Co, Mg

↑ ↓ ↑ ↓ ↑ ↓ ↑ ↓

Two sublattices A and B.



$MgAl_2O_4$

Spinel crystal structure XY_2O_4

8 tetrahedral sites A (surrounded by 4 O) $5\mu_B$ ↑

16 octahedral sites B (surrounded by 6 O) $9\mu_B$ ↓

per unit cell

Ferrimagnets

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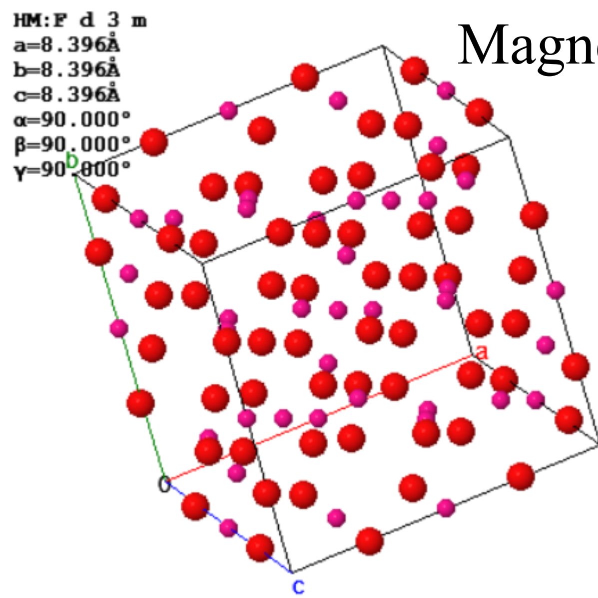
Ferrites $MO \cdot Fe_2O_3$

M = Fe, Zn, Cd, Ni, Cu, Co, Mg

$\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$



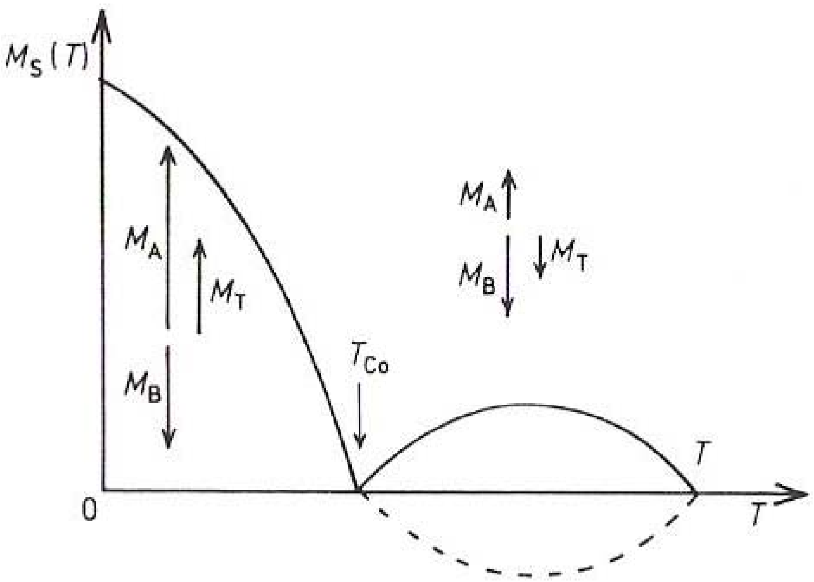
Magnetite

Exchange integrals J_{AA}, J_{AB}, and J_{BB} are all negative (antiparallel preferred)

$|J_{AB}| > |J_{AA}|, |J_{BB}|$

↑↓↑↓↑↓↑↓ Ferrimagnetism

gauss = 10^{-4} T
oersted = $10^{-4}/4\pi\times10^{-7}$ A/m



D. Gignoux, magnetic properties of Metallic systems

Table 33.3
SELECTED FERRIMAGNETS, WITH CRITICAL TEMPERATURES T_c AND SATURATION MAGNETIZATION M_0

MATERIAL	T_c (K)	M_0 (gauss)[a]
Fe_3O_4 (magnetite)	858	510
$CoFe_2O_4$	793	475
$NiFe_2O_4$	858	300
$CuFe_2O_4$	728	160
$MnFe_2O_4$	573	560
$Y_3Fe_5O_{12}$ (YIG)	560	195

[a] At $T = 0$(K).
Source: F. Keffer, *Handbuch der Physik*, vol. 18, pt. 2, Springer, New York, 1966.

Kittel

Antiferromagnetism

Negative exchange energy $J_{AB} < 0$.

$$\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow$$

At low temperatures, below the Neel temperature T_N, the spins are aligned antiparallel and the macroscopic magnetization is zero.

Spin ordering can be observed by neutron scattering.

At high temperature antiferromagnets become paramagnetic. The macroscopic magnetization is zero and the spins are disordered in zero field.

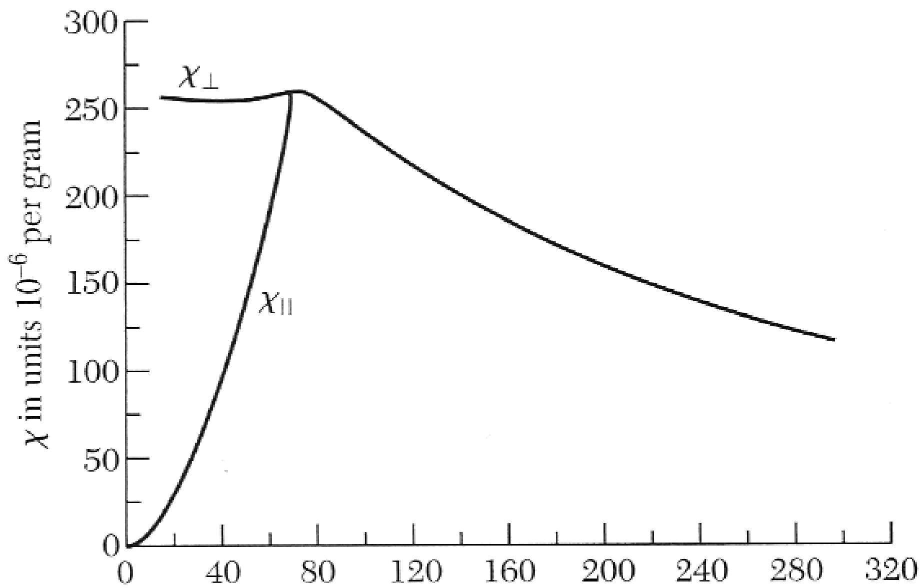
$$\chi = \mu_0 \frac{\vec{M}_A + \vec{M}_B}{\vec{B}_a} = \frac{C}{T + \Theta}$$

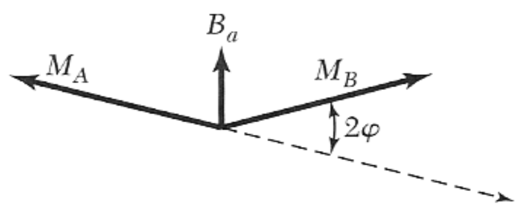
Curie-Weiss temperature

Antiferromagnetism ↑↓↑↓↑↓↑↓

Average spontaneous magnetization is zero at all temperatures.

Source: Kittel





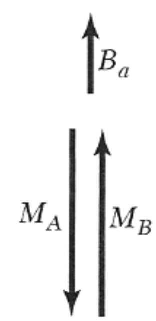
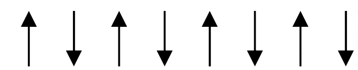
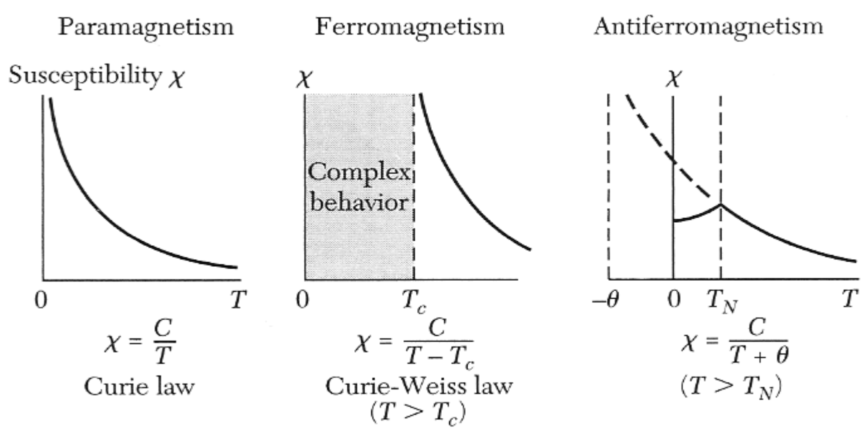


Table 2 Antiferromagnetic crystals



Substance	Paramagnetic ion lattice	Transition temperature, T_N, in K	Curie-Weiss θ, in K	$\frac{\theta}{T_N}$	$\frac{\chi(0)}{\chi(T_N)}$
MnO	fcc	116	610	5.3	$\frac{2}{3}$
MnS	fcc	160	528	3.3	0.82
MnTe	hex. layer	307	690	2.25	
MnF_2	bc tetr.	67	82	1.24	0.76
FeF_2	bc tetr.	79	117	1.48	0.72
$FeCl_2$	hex. layer	24	48	2.0	<0.2
FeO	fcc	198	570	2.9	0.8
$CoCl_2$	hex. layer	25	38.1	1.53	
CoO	fcc	291	330	1.14	
$NiCl_2$	hex. layer	50	68.2	1.37	
NiO	fcc	525	~2000	~4	
Cr	bcc	308			



from Kittel

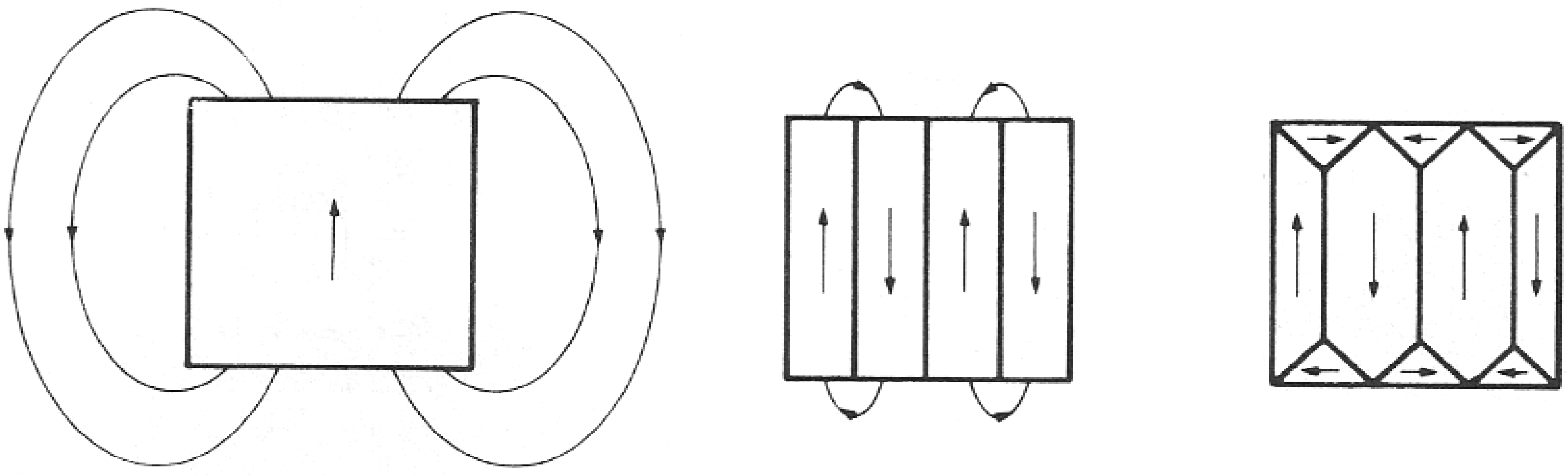
Applications of magnetism

Hard magnets: permanent magnets, motors, generators, microphones

Soft magnets: transformers

Magnetic recording

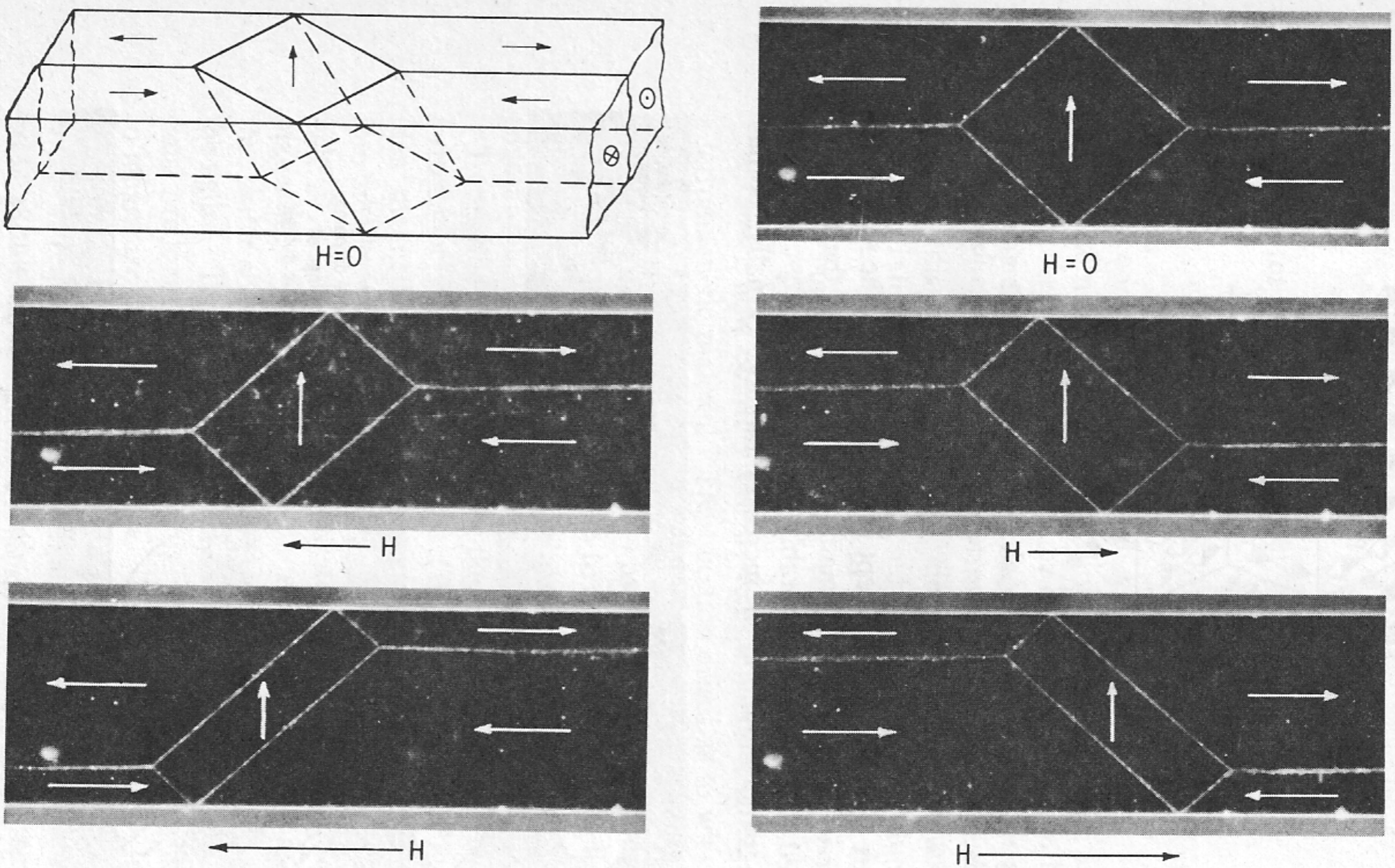
Magnetic domains (weisssche Bezirke)



Magnetic energy density

$$\frac{B^2}{2\mu_0}$$

Costs energy to introduce domain walls where spin up regions are adjacent to spin down regions.



H=0
H=0
H
H
H
H

Ferromagnetic domains

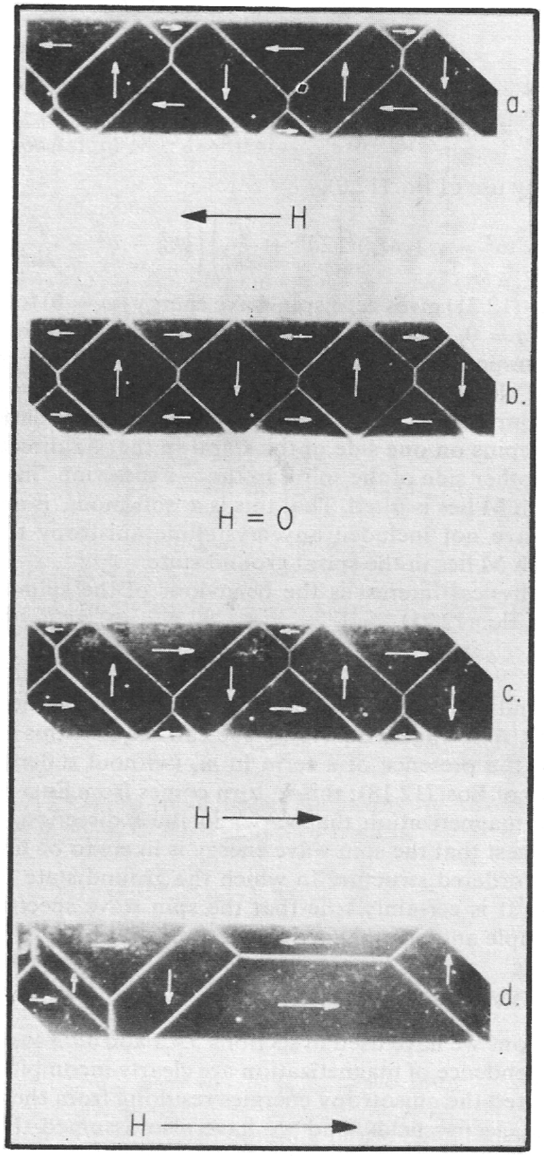
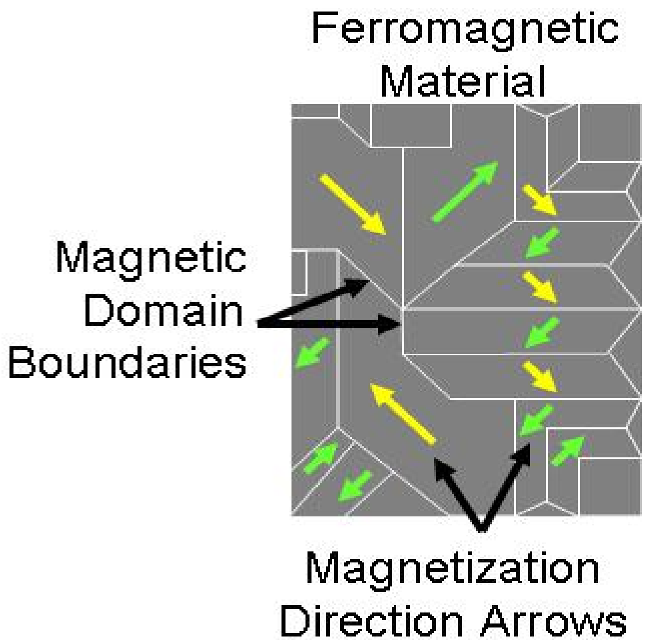


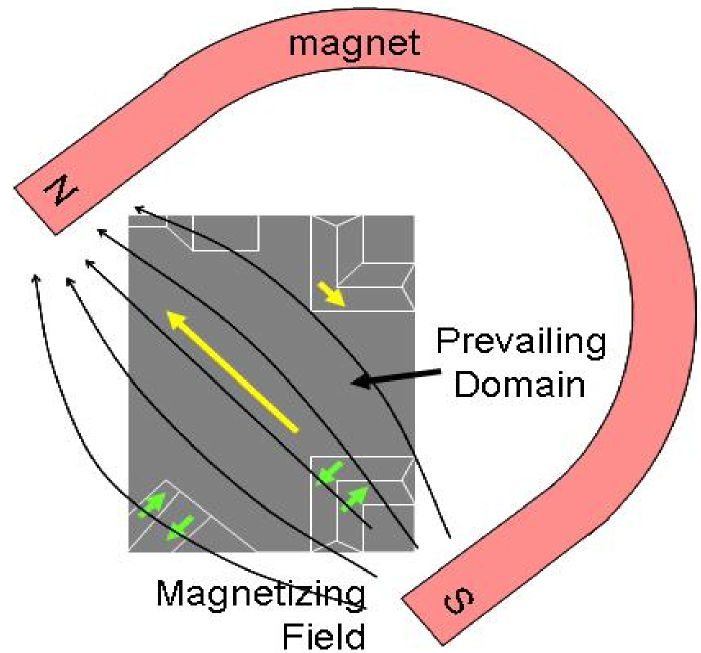
Fig. 12.5. Photographs showing reversible domain wall motion in a 50 μm whisker from (*a*) to (*b*) to (*c*), with an irreversible jump from (*c*) to (*d*).
{R. W. de Blois and C. D. Graham, *J. Appl. Phys.*, **29**, 931 (1958)}.

Weak fields: favorable domains expand
Strong fields: domains rotate to align with field

Irreproducible jump between c and d.

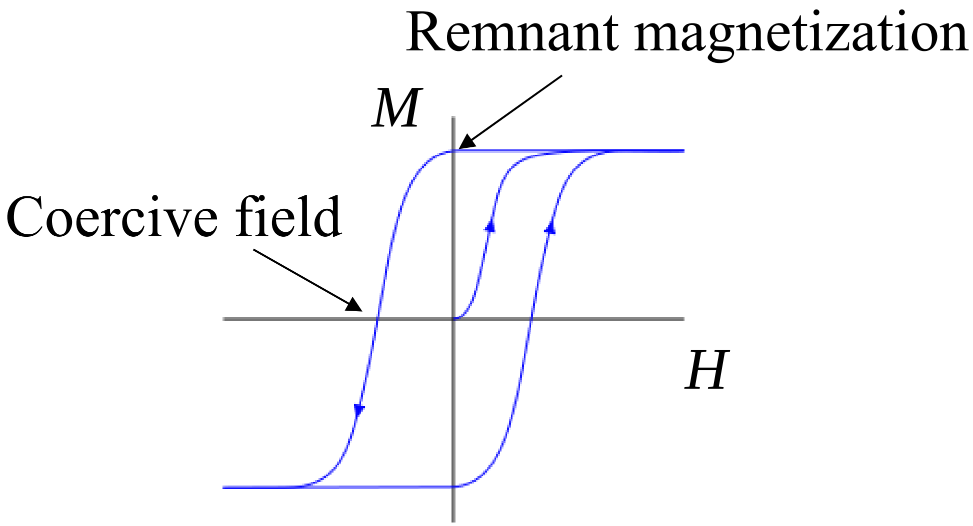
Magnetizing a magnet





Weak fields: favorable domains expand
Strong fields: domains rotate to align with field

Hysteresis



$$B = \mu_0 \left(H + M \right)$$

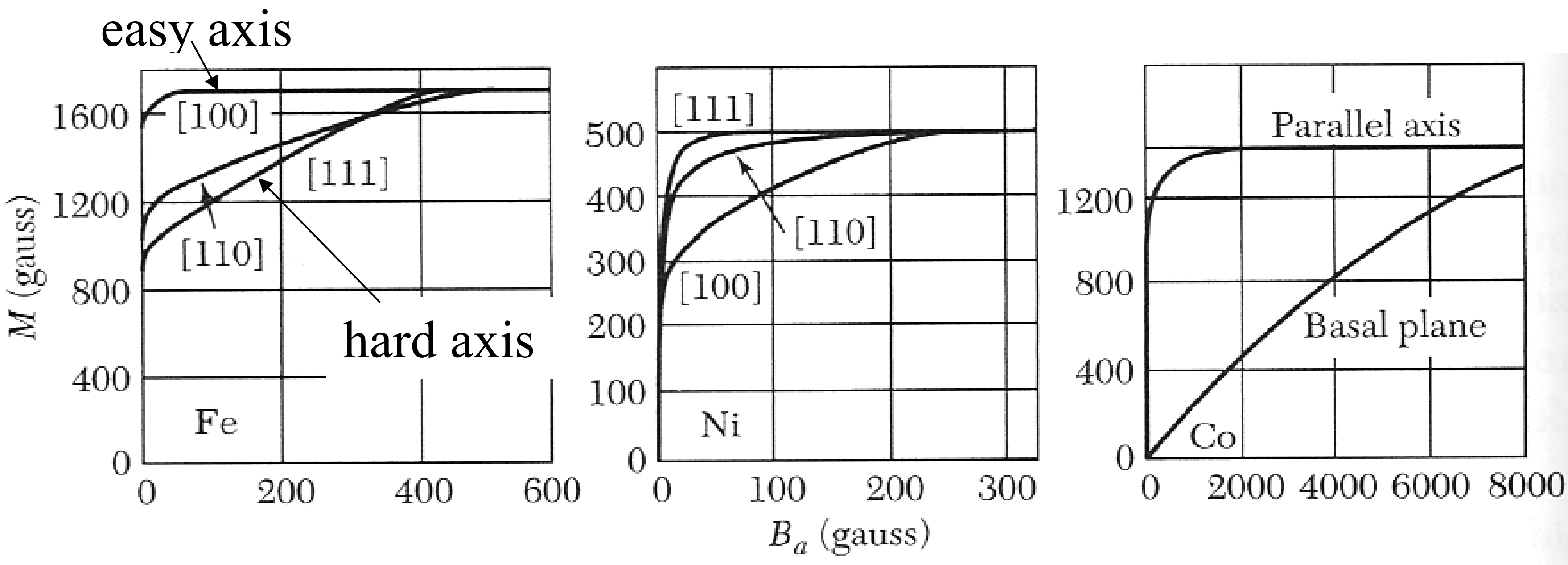
$$M = \chi H$$

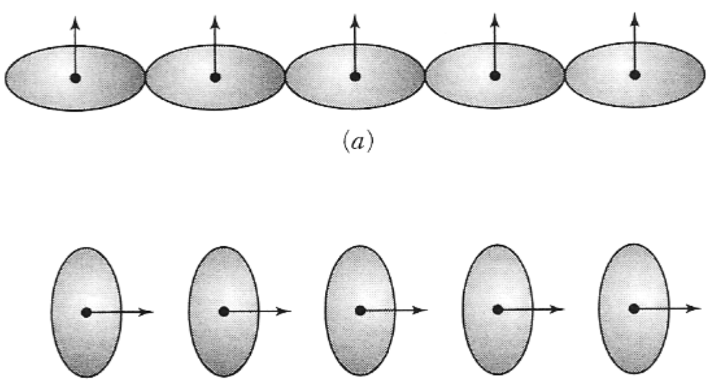
$$B = \mu_0 (1 + \chi) H = \mu_r \mu_0 H$$

$$\mu_r = \left(1 + \chi \right)$$

Area of the loop is proportional to energy dissipated in traversing the loop.

Anisotropy energy




(a)

Spin-orbit coupling couples the shape of the wavefunction to the spin. The exchange energy depends on the overlap of the wavefunctions and thus on spin direction.

Bloch wall

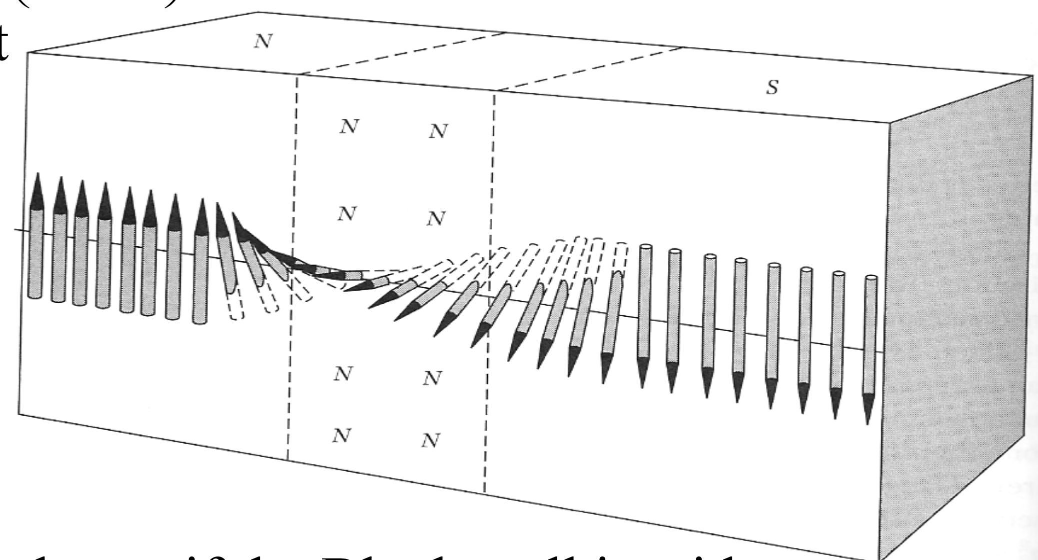
energy of two spins

$$w = -J\vec{S}_i \cdot \vec{S}_j = -JS^2 \cos\varphi \approx -JS^2\left(1 - \frac{\varphi^2}{2}\right)$$

neglecting the constant part

$$w \approx JS^2 \frac{\varphi^2}{2}$$

$$\varphi = \frac{\pi}{N}$$



energy of a line of spins

$$Nw \approx \frac{JS^2\pi^2}{2N}$$

energy is lower if the Bloch wall is wide

energy per unit area $\approx \dfrac{JS^2\pi^2}{2Na^2}$

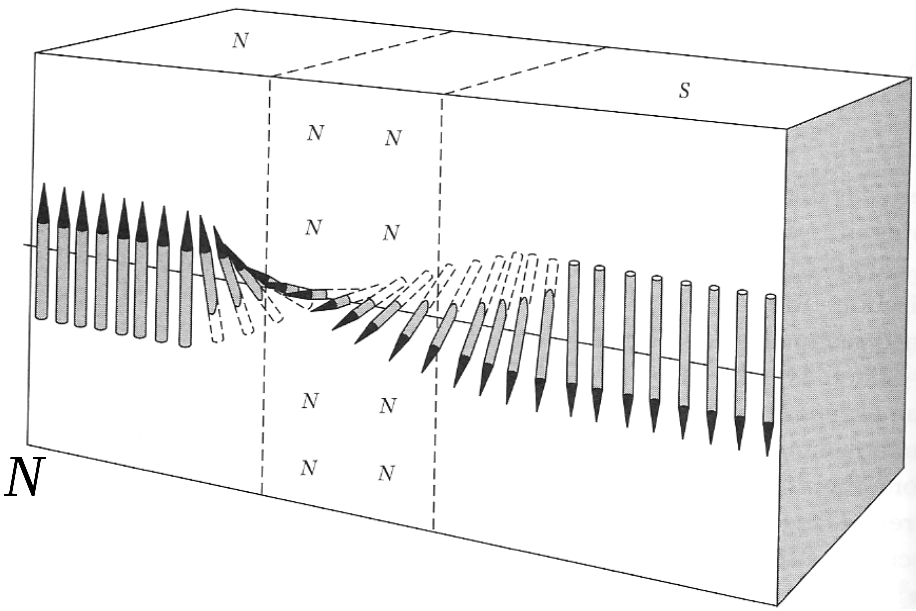
a is the lattice constant

Bloch wall

Anisotropy energy depends on the number of spins pointing in the hard direction

$\approx KNa$ ← Na = thickness of wall

↑ anisotropy constant J/m^3



Total energy per unit area:

$$E = \frac{JS^2\pi^2}{2Na^2} + KNa \qquad [\mathrm{J/m^2}]$$

smaller for large N smaller for small N

$$\frac{dE}{dN} = 0 \Rightarrow -\frac{JS^2\pi^2}{2N^2a^2} + Ka = 0$$

$$N = \sqrt{\frac{JS^2\pi^2}{2Ka^3}}$$

$N \sim 300$ for iron