

Advanced Solid State Physics

Optical properties of insulators and semiconductors

In an insulator, all charges are bound. By applying an electric field, the electrons and ions can be pulled out of their equilibrium positions. When this electric field is turned off, the charges oscillate as they return to their equilibrium positions. A simple model for an insulator can be constructed by describing the motion of the charge as a damped mass-spring system. The differential equation that describes the motion of a charge is,

$$m \, \frac{d^2 x}{dt^2} + b \, \frac{dx}{dt} + kx = -qE.$$

Rewriting above equation using $\omega_0=\sqrt{\frac{k}{m}}$ and the damping constant $\gamma=\frac{b}{m}$ yields,

 $rac{d^{\,2}x}{dt^{\,2}}+\gammarac{dx}{dt}+\omega_{0}^{\,2}x=-rac{qE}{m}\,.$

If the electric field is pulsed on, the response of the charges is described by the impulse response function g(t). The impulse response function satisfies the equation,

$$rac{d^2g}{dt^2}+\gammarac{dg}{dt}+\omega_0^2g=-rac{q}{m}\,\delta(t).$$

The solution to this equation is zero before the electric field is pulsed on and at the time of the pulse the charges suddenly start oscillating with the frequency $\omega_1 = \sqrt{\omega_0^2 - \frac{\gamma^2}{4}}$. The amplitude of the oscillation decays exponentially to zero in a characteristic time $\frac{2}{\gamma}$.

$$g(t) = -\frac{q}{m\omega_1} \exp(-\frac{\gamma}{2} t) \sin(\omega_1 t).$$



Outline Quantization Photons Electrons Magnetic effects and Fermi surfaces Linear response Transport **Crystal Physics** Electron-electron interactions Quasiparticles Structural phase transitions Landau theory of second order phase transitions Superconductivity Exam guestions Appendices Lectures Books Course notes TUG students Making presentations

Dielectrics

Dielectrics used as electrical insulators should not conduct.

Large breakdown field.

Low AC losses.

Sometimes a low dielectric constant is desired (CMOS interconnects)

Sometimes a high dielectric constant is desired (supercapacitors).



Electrode separation in meters $\times \ 10^{-2}$

AC losses - loss tangent

In an ideal capacitor, current leads voltage by 90°.

Because the dielectric constant is complex, in real materials current leads voltage by 90° - δ .

Power loss =
$$\frac{\omega \varepsilon_1 V_0^2}{2} \tan \delta$$

Becomes more of an issue at high frequencies (microwaves)



Loss tangent

Substance	Dielectric Constant (relative to air)	Dielectric Strength (Y/mil)	Loss Tangent	Max Temp (°F)	
ABS (plastic), Molded	2.0 - 3.5	400 - 1350	0.00500 - 0.0190	171 - 228	
Air	1.00054	30 - 70			
Alumina - 96% - 99.5%	10.0 9.6		0.0002 @ 1 GHz 0.0002 @ 100 MHz 0.0003 @ 10 GHz		
Aluminum Silicate	5.3 - 5.5				
Bakelite	3.7				
Bakelite (mica filled)	4.7	325 - 375			
Balsa Wood	1.37 @ 1 MHz 1.22 @ 3 GHz		0.012 @ 1 MHz 0.100 @ 3 GHz		
Beeswax (yellow)	2.53 @ 1 MHz 2.39 @ 3 GHz		0.0092 @ 1 MHz 0.0075 @ 3 GHz		
Beryllium oxide	6.7		0.006 @ 10 GHz		
Butyl Rubber	2.35 @ 1 MHz 2.35 @ 3 GHz		0.001 @ 1 MHz 0.0009 @ 3 GHz		
Carbon Tetrachloride	2.17 @ 1 MHz 2.17 @ 3 GHz		<0.0004 @ 1 MHz 0.0004 @ 3 GHz		
Diamond	5.5 - 10				
Delrin (acetyl resin)	3.7	500		180	
Douglas Fir	1.9 @ 1 MHz		0.023 @ 1 MHz		
Douglas Fir Plywood	1.93 @ 1 MHz 1.82 @ 3 GHz		0.026 @ 1 MHz 0.027 @ 3 GHz		
Enamel	5.1	450			
Epoxy glass PCB	5.2	700			
Ethyl Alcohol (absolute)	24.5 @ 1 MHz 6.5 @ 3 GHz		0.09 @ 1 MHz 0.25 @ 3 GHz		
Ethylene Glycol	41 @ 1 MHz 12 @ 3 GHz		-0.03 @ 1 MHz 1 @ 3 GHz		
Formica XX	4.00				
FR-4 (G-10) - low resin	4.9		0.008 @ 100 MHz		
- high resin	4.2		0.008 @ 3 GHz		
Fused quartz	3.8		0.0002 @ 100 MHz 0.00006 @ 3 GHz		
Fused silica (glass)	3.8				
Gallium Arsenide (GaAs)	13.1		0.0016 @ 10 GHz		
Germanium	16				
Glass	4 - 10				
Glass (Corning 7059)	5.75		0.0036 @ 10 GHz		
Gutta-percha	2.6				
Halowax oil	4.8				
High Density Polyethylene (HDPE), Molded	1.0 - 5.0	475 - 3810	0.0000400 - 0.00100	158 - 248	
Ice (pure distilled water)	4.15 @ 1 MHz 3.2 @ 3 GHz		0.12 @ 1 MHz 0.0009 @ 3 GHz		
Kapton® Type 100 Type 150	3.9 2.9	7400 4400		500	

Polarizability

Overdamped modes

- Orientation polarizability
- Space charge polarizability

Underdamped modes

- Ionic polarizability
- Electronic polarizability

Orientation (dipolar) Polarizability

For materials (gases, liquids, solids) with a permanent dipole moment.

The theory is very similar to paramagnetism.



$$\chi \propto \frac{1}{T}$$
 Curie law

Orientation Polarizability



Orientation (dipolar) Polarizability



For low frequencies the dipoles can reorient with the field but at high frequencies they can't respond fast enough.

Water



Space charge polarizability

Multiple phases are present where one phase has a much higher resistivity than the other. Charge accumulates at the interfaces of the phases.

Like a network of resistors and capacitors. This results in an overdamped mode.







Ionic Polarizability

Displacement of ions of opposite sign. Only in ionic substances.



This is an underdamped mode in the infrared.

Electronic polarizability (all materials)



-

$$\vec{P} = N\vec{p} = N\alpha\vec{E}$$

density polarizability

Table 1 Electronic polarizabilities of atoms and ions, in 10⁻²⁴ cm³

			He	Li ⁺	Be^{2+}	B^{3+}	C^{4+}
Pauling JS			0.201	$0.029 \\ 0.029$	0.008	0.003	0.0013
Pauling JS-(TKS)	O^{2-} 3.88 (2.4)	${ m F}^-$ 1.04 0.858	Ne 0.390	Na ⁺ 0.179 0.290	${ m Mg}^{2+} \ 0.094$	Al^{3+} 0.052	Si ⁴⁺ 0.0165
Pauling JS-(TKS)	${S^{2-}}\ 10.2\ (5.5)$	Cl⁻ 3.66 2.947	Ar 1.62	${ m K}^+ \ 0.83 \ 1.133$	${ m Ca}^{2+}$ 0.47 (1.1)	$\frac{\mathrm{Se}^{3+}}{0.286}$	${ m Ti}^{4+} \ 0.185 \ (0.19)$
Pauling JS-(TKS)	${ m Se}^{2-}\ 10.5\ (7.)$	Br ⁻ 4.77 4.091	Kr 2.46	${ m Rb}^+ \ 1.40 \ 1.679$	${ m Sr}^{2+} \ 0.86 \ (1.6)$	Y ³⁺ 0.55	$ m Zr^{4+}$ 0.37
Pauling JS-(TKS)	Te^{2-} 14.0 (9.)	I ⁻ 7.10 6.116	Xe 3.99	Cs ⁺ 2.42 2.743	${ m Ba}^{2+}\ 1.55\ (2.5)$	La ³⁺ 1.04	Ce^{4+} 0.73

Polarizability



Inter- and intraband transitions

When the bands are parallel, there is a peak in the absorption (ε ")



Optical spectroscopy has developed into the most important experimental tool for band structure determination. - Kittel