$$M(k_{x},k_{y},k_{z}) \coloneqq \begin{bmatrix} 4 - \cos\left[\left(k_{x}+k_{y}\right)\cdot\frac{a_{0}}{2}\right] - \cos\left[\left(k_{x}-k_{y}\right)\cdot\frac{a_{0}}{2}\right] - \cos\left[\left(k_{x}-k_{z}\right)\cdot\frac{a_{0}}{2}\right] \\ -\cos\left[\left(k_{x}-k_{z}\right)\cdot\frac{a_{0}}{2}\right] + \cos\left[\left(k_{x}-k_{z}\right)\cdot\frac{a_{0}}{2}\right] \\ -\cos\left[\left(k_{x}-k_{z}\right)\cdot\frac{a_{0}}{2}\right] \\ -\cos\left[\left(k_{x}-k$$

 $\lambda \left(k_{x}^{},k_{y}^{},k_{z}^{} \right) \coloneqq \text{eigenwerte} \left(M \left(k_{x}^{},k_{y}^{},k_{z}^{} \right) \right)$







% solid state physics % calculation of dispersion relations for monoatomic 3D lattices % Philipp Thaler % 4.5.09 % parameters a = 4.49; % length of a conventional unit cell in A % spring force constant in N/m C = 4/2; $m = 0.344 \times 10^{-24};$ % mass in kg % postitions of the examined atom and the nearest neighbors % fcc: an = a/2 * [0,0,0; 1,1,0; 1,0,1; 0,1,1; -1,-1,0; -1,0,-1; 0,-1,-1; 1,-1,0; 1,0,-1; 0,1,-1; -1,1,0; -1,0,1; 0,1,-1; -1,1,0; -1,0,1; 0,-1,1]; % wavevector to plot 100-110-111 for fcc nk = 100;k0 = linspace(0, 0, nk);k100 = linspace(0, 2*pi, nk)/a; k110 = linspace(-2*pi*sqrt(2),0,nk)/a; k111 = linspace(0, pi*sqrt(3), nk)/a; kx = [k100, k110/sqrt(2), k111/s k111/sqrt(3)]; k111/sqrt(3)]; ky = [k0,k110/sqrt(2), kz = [k0]k111/sqrt(3)] k0, k100(end)+fliplr(-k110), k100(end)-k110(1)+k111]; kk = [k100,% calculate the frequencies w = NaN(3, length(kk)); for n = 1:length(kk) k = [kx(n), ky(n), kz(n)];w(:,n) = disp_solver(an, m, C, k); end %plot fcc: figure(1) w(1,:), 'b-', kk, w(2,:), 'b-', kk, w(3,:), 'b-') pl ot (kk, hold on %Points Gamma, X, K, Gamma, L ymax = round(10*1.2*max(w(:)))/10; plot([0,0],[0,ymax],'r-') plot([2*pi,2*pi]/a,[0,ymax],'r-') plot([2*pi+2*pi*sqrt(2)-sqrt(9/2)*pi,2*pi+2*pi*sqrt(2)-sqrt(9/2)*pi]/a,[0,ymax],'r-') plot([2*pi+2*pi*sqrt(2),2*pi+2*pi*sqrt(2)]/a,[0,ymax],'r-') plot([2*pi+2*pi*sqrt(2)+pi*sqrt(3),2*pi+2*pi*sqrt(2)+pi*sqrt(3)]/a,[0,ymax],'r-') xlim([kk(1) kk(end)])
ylim([0 ymax])
hold off

```
solver for the eigenvalue problem of 3d dispersion relations
%
   philipp thaler
%
   8.5.09
function w = disp solver(an, m, C, k)
%%%%%%%%%%%
% solver %
%%%%%%%%%%%
% directions of the neighbors
ar =an(2: end, :);
arnorm = zeros(si ze(ar, 1), 1);
for n = 1: si ze(ar, 1)
       arnorm(n) = norm(ar(n, [1, 2, 3]));
       ar(n, [1, 2, 3]) = ar(n, [1, 2, 3]). /arnorm(n);
end
% Newton's law
% m * d2u/dt2 = sum( C * delta_u)
% approach
\% u = u0 * exp(-iwt) * exp(iak)
% split in 3 dimensions, using the approach for u
   -m/C * w^2 * ux + sum(del ta_u) = 0
-m/C * w^2 * uy + sum(del ta_u) = 0
%
%
   -m/C * w^2 * uz + sum(delta_u) = 0
%
% delta_u = u(k',l',m')-u(k,l,m)
% for nearest neighbor #1 k'=k+1, l'=l+1, m'=m
% --> delta_u = u0 * exp(-iwt) * exp(i(k*kx+l*ky+m*kz)) * (exp((k'-k)kx + (l'-l)ky +
(m' - m)kz)) - 1)
for n = 1: size(ar, 1)
       % x-direction, sum(delta_u):
       del ta_uxx = (ar(:,1).*ar(:,1)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uxy = (ar(:,1).*ar(:,2)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uxz = (ar(:,1).*ar(:,3)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
       % y-direction, sum(del ta_u):
del ta_uyx = (ar(:,2).*ar(:,1)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uyy = (ar(:,2).*ar(:,2)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uyz = (ar(:,2).*ar(:,3)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
       % z-direction, sum(delta_u):
       del ta_uzx = (ar(:,3).*ar(:,1)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uzy = (ar(:,3).*ar(:,2)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
del ta_uzz = (ar(:,3).*ar(:,3)).*(exp(i*sum(an(2:end,:).*repmat(k,size(ar,1),1),2))-1);
end
```

M =
-real ([sum(delta_uxx, 1), sum(delta_uxy, 1), sum(delta_uxz, 1); sum(delta_uyx, 1), sum(delta_uyy, 1),
sum(delta_uyz, 1); sum(delta_uzx, 1), sum(delta_uzy, 1), sum(delta_uzz, 1)]);
% eigs(M) = w^2 * m/C
w = real(sqrt(eigs(M)*C/m));