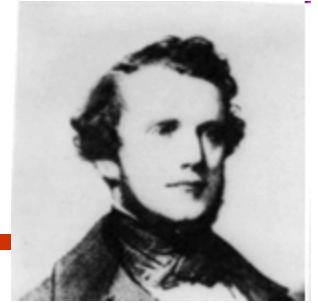


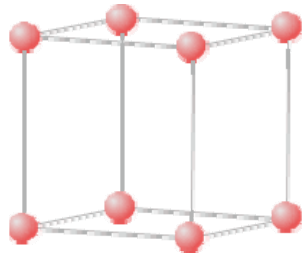
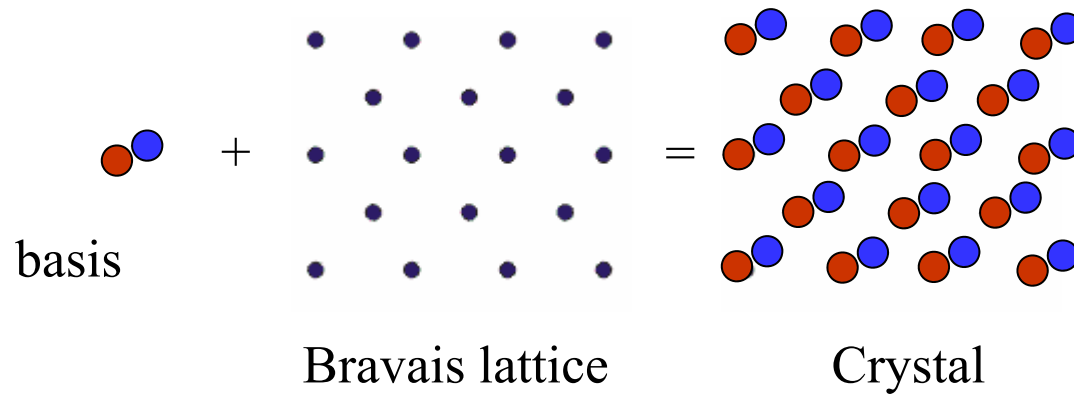
7. Crystal structure, Diffraction

April 10, 2018

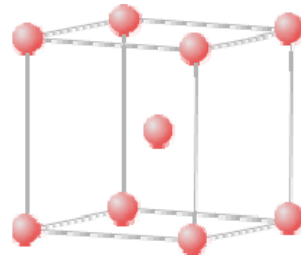
Bravais lattice



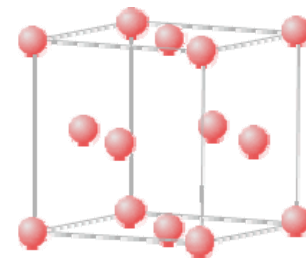
Auguste Bravais



simple cubic



body centered
cubic, bcc



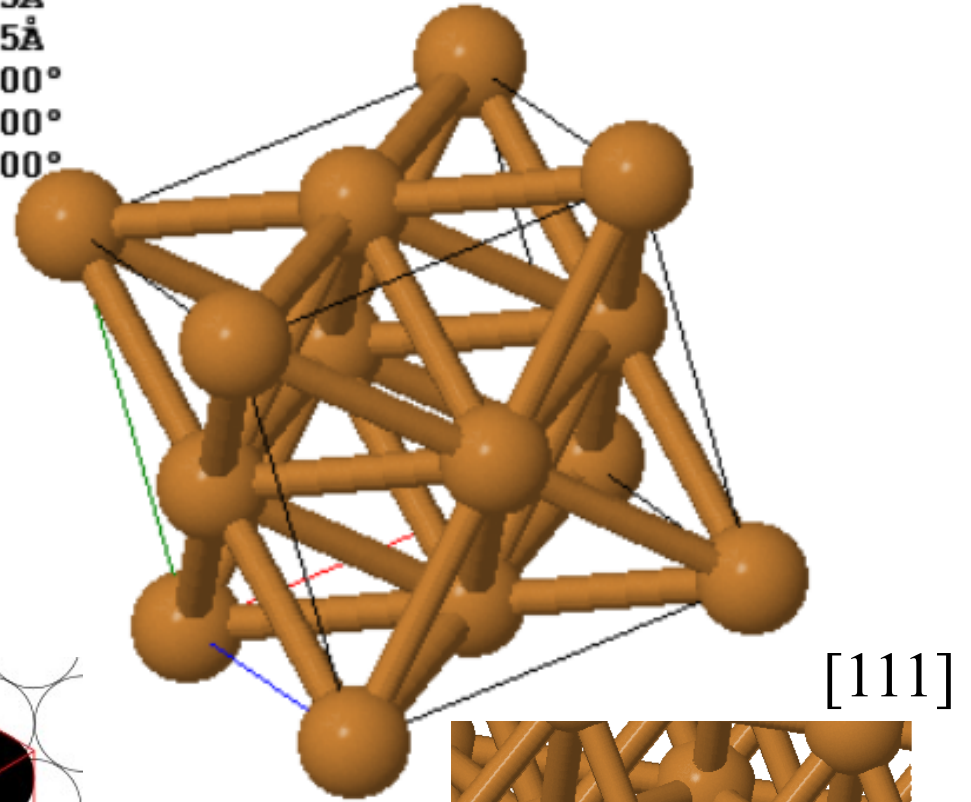
face centered
cubic, fcc

fcc

Number 225

Al, Cu,
Ni, Sr,
Rh, Pd,
Ag, Ce,
Tb, Ir,
Pt, Au,
Pb, Th

HM: $F m \bar{3} m$
 $a=3.615\text{\AA}$
 $b=3.615\text{\AA}$
 $c=3.615\text{\AA}$
 $\alpha=90.000^\circ$
 $\beta=90.000^\circ$
 $\gamma=90.000^\circ$

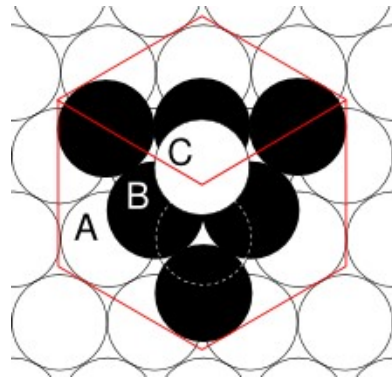


Primitive Vectors:

$$\vec{a}_1 = \frac{a}{2} \hat{y} + \frac{a}{2} \hat{z}$$

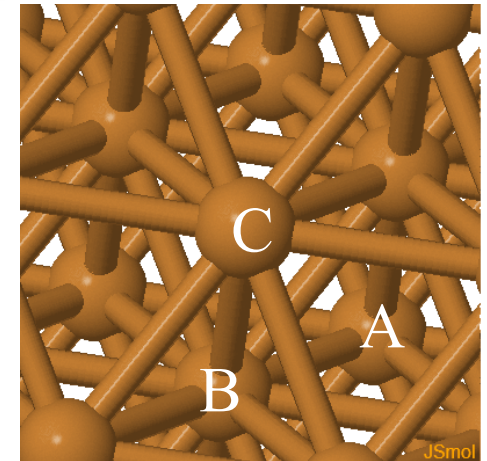
$$\vec{a}_2 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{z}$$

$$\vec{a}_3 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{y}$$



Basis Vector:

$$\vec{B}_1 = (0, 0, 0)$$



bcc

W Number 229

Na

K

V

Cr

Fe

Rb

Nb

Mo

Cs

Ba

Eu

Ta

Primitive Vectors:

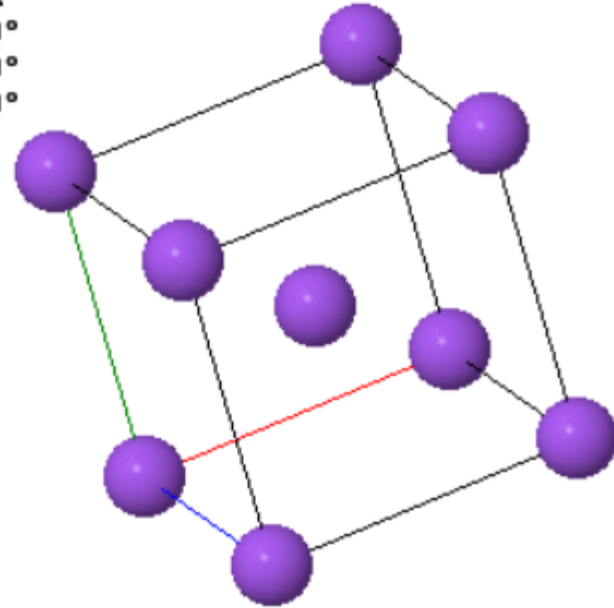
$$\vec{a}_1 = -\frac{a}{2}\hat{x} + \frac{a}{2}\hat{y} + \frac{a}{2}\hat{z}$$

$$\vec{a}_2 = \frac{a}{2}\hat{x} - \frac{a}{2}\hat{y} + \frac{a}{2}\hat{z}$$

$$\vec{a}_3 = \frac{a}{2}\hat{x} + \frac{a}{2}\hat{y} - \frac{a}{2}\hat{z}$$

Basis Vector: $\vec{B}_1 = (0, 0, 0)$

HM: 1M-3M
a=4.291Å
b=4.291Å
c=4.291Å
α=90.000°
β=90.000°
γ=90.000°

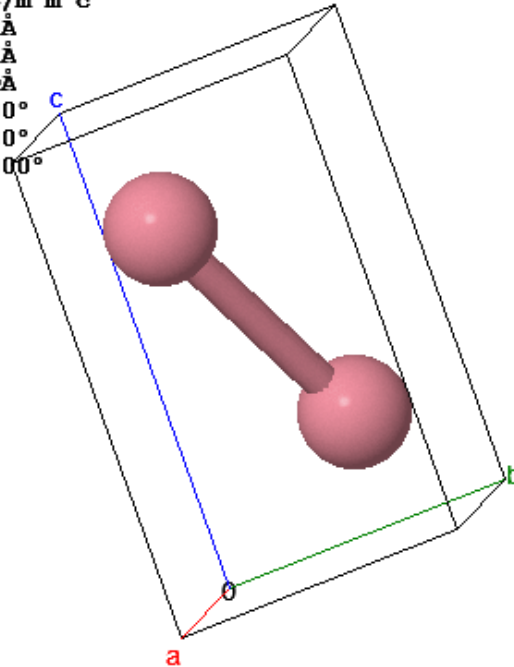


hcp

Space group 194 Crystallographic unit cell

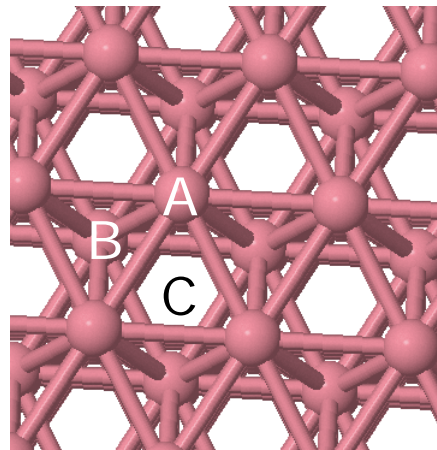
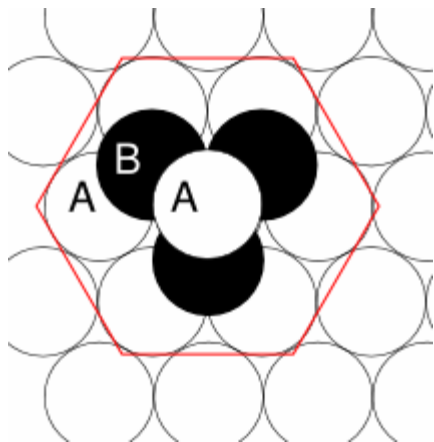
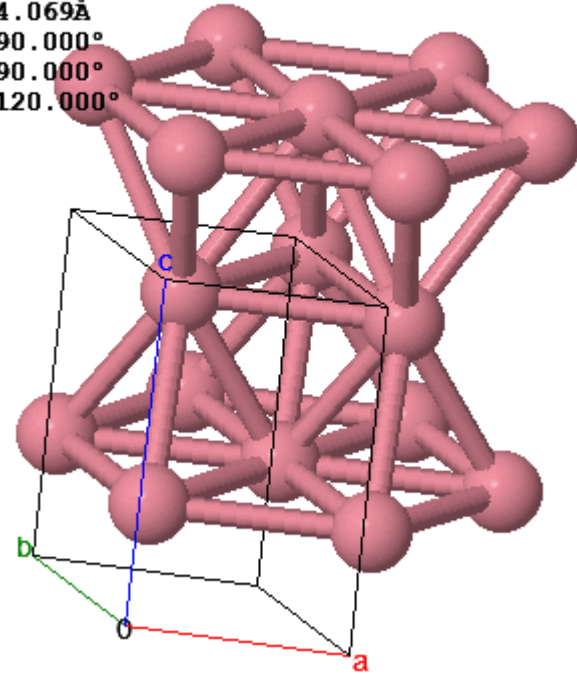
Mg, Be,
Sc, Ti,
Co, Zn,
Y, Zr, Tc,
Ru, Cd,
Gd, Tb,
Dy, Ho,
Er, Tm,
Lu, Hf,
Re, Os,
Tl

HM: P 63/m m c
a=2.507Å
b=2.507Å
c=4.069Å
α=90.000°
β=90.000°
γ=120.000°



Hexagonal unit cell

HM: P 63/m m c
a=2.507Å
b=2.507Å
c=4.069Å
α=90.000°
β=90.000°
γ=120.000°

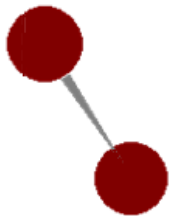


Hexagonal Bravais lattice
Basis vectors:

$$\vec{B}_1 = (0, 0, 0) \quad \vec{B}_2 = \left(\frac{2}{3}, \frac{1}{3}, \frac{1}{2}\right)$$

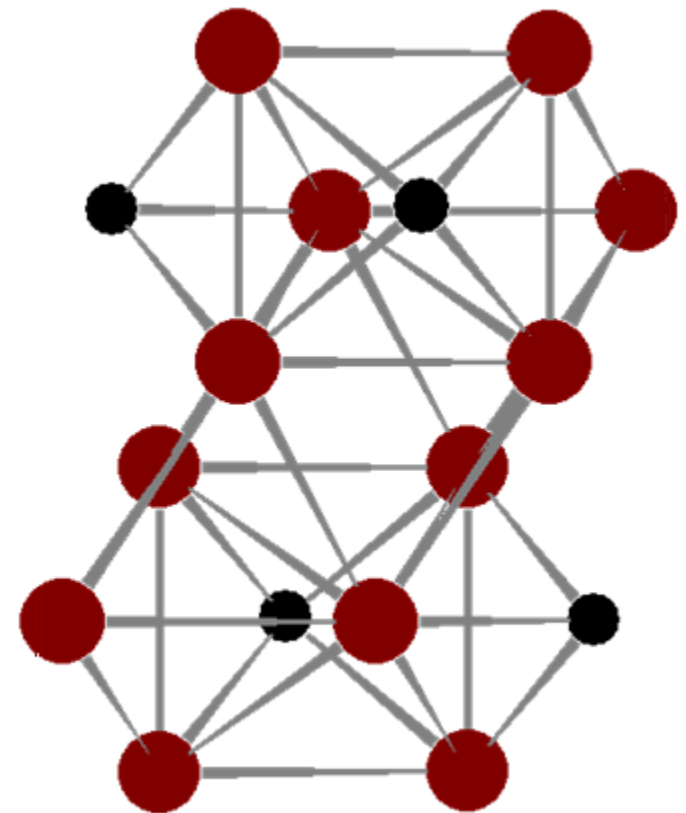
Asymmetric unit

```
cell 5.09000 6.74800 4.52300 90.000 90.000 90.000
natom 3
Fe1 26 0.18600 0.06300 0.32800
Fe2 26 0.03600 0.25000 0.85200
C 6 0.89000 0.25000 0.45000
rgnr 62
Cohenite (Cementite) Fe3 C
```

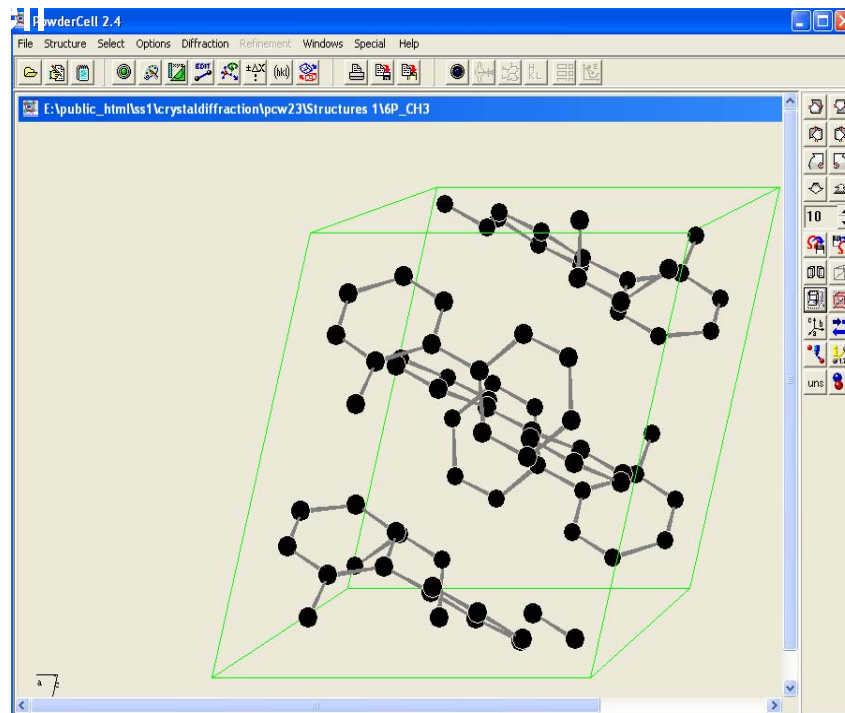


Asymmetric unit

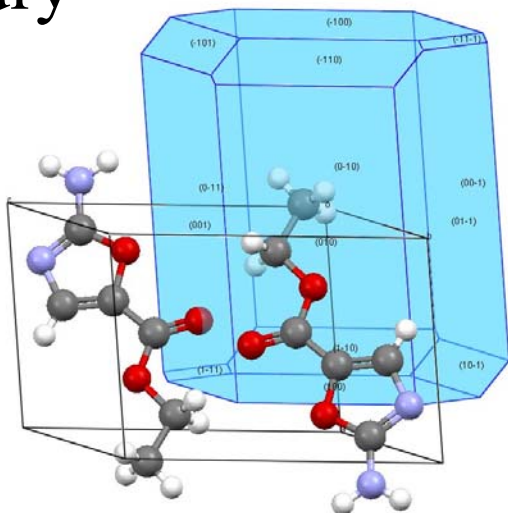
Unit cell



PowderCell



Mercury

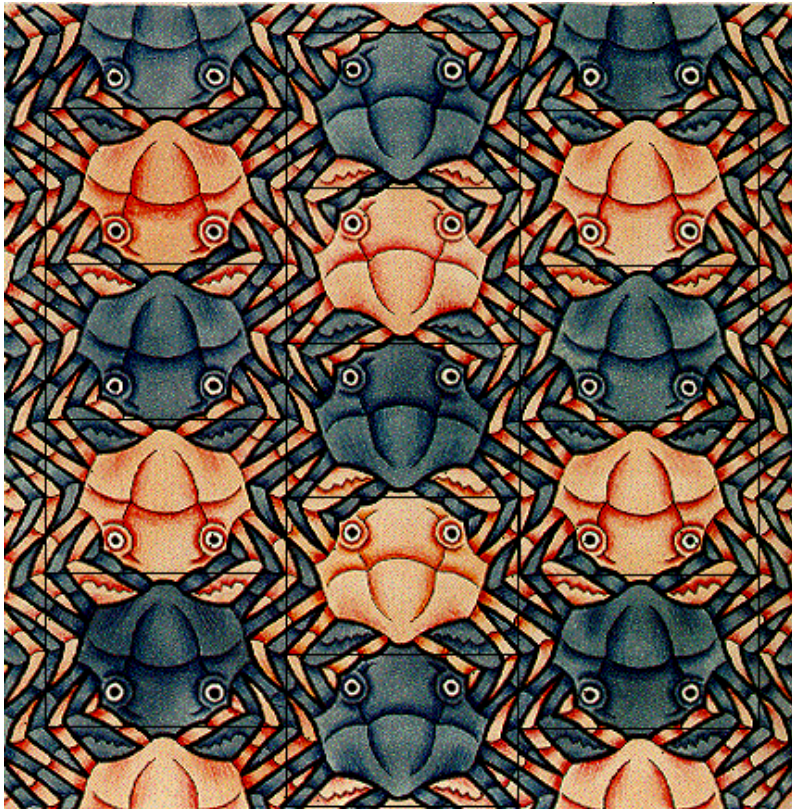


Standard data file: *.cif

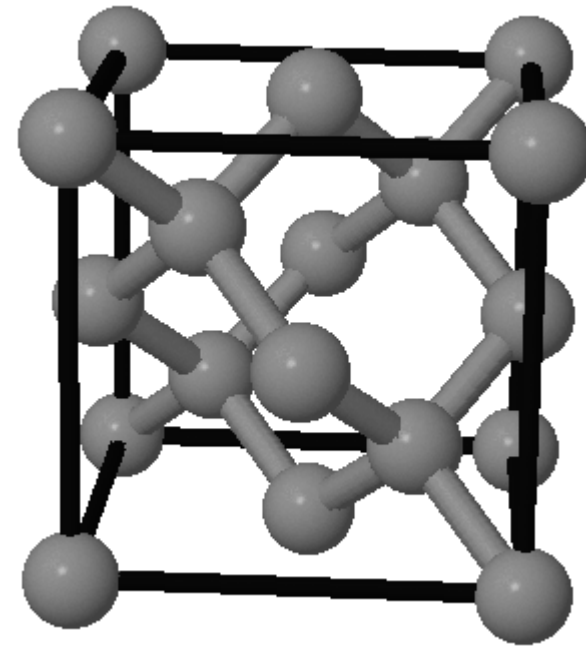
http://www.bam.de/de/service/publikationen/powder_cell.htm

<https://www.ccdc.cam.ac.uk/Community/csd-community/freemercury/>

Inequivalent atoms in the unit cell

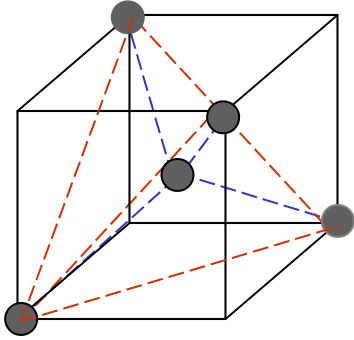


An element can have two distinct positions



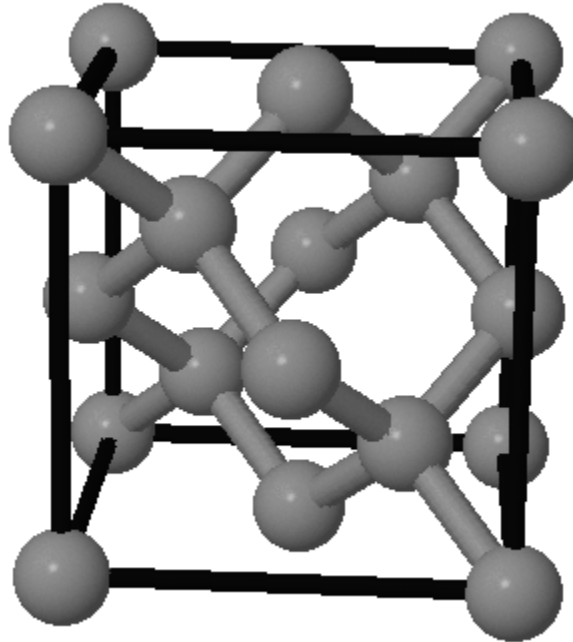
Diamond conventional unit cell

Diamond



$$a = b = c, \quad \alpha = 90^\circ, \beta = 90^\circ, \gamma = 90^\circ$$

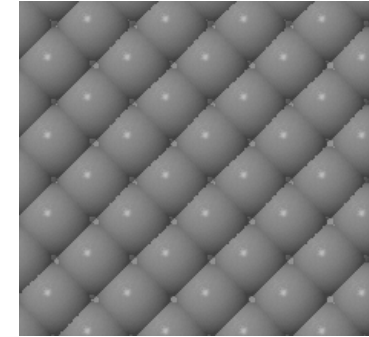
Space group: 227
point group: m3m



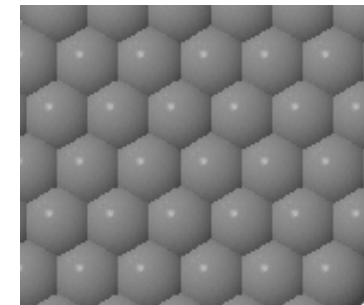
Primitive lattice vectors:

$$\vec{a}_1 = \frac{a}{2}\hat{x} + \frac{a}{2}\hat{y}, \quad \vec{a}_2 = \frac{a}{2}\hat{x} + \frac{a}{2}\hat{z}, \quad \vec{a}_3 = \frac{a}{2}\hat{y} + \frac{a}{2}\hat{z}.$$

Basis: $\vec{B}_1 = (0, 0, 0), \quad \vec{B}_2 = (0.25, 0.25, 0.25).$



$$(100): \frac{2}{a^2}$$



$$(111): \frac{4}{\sqrt{3}a^2}$$

zincblende

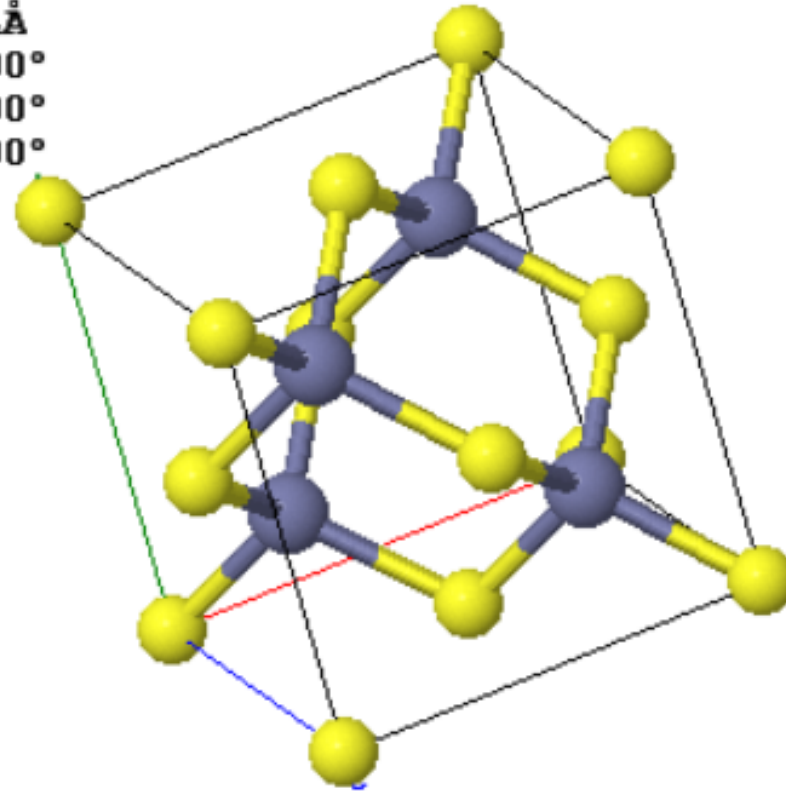
ZnS

GaAs

InP

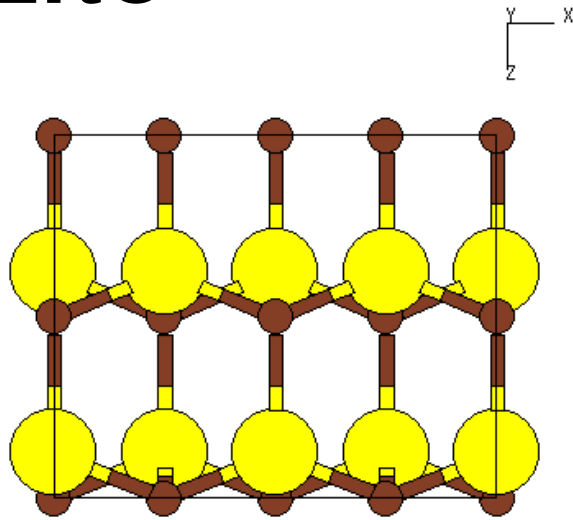
HM: $F\bar{4}3M$
 $a=5.434\text{\AA}$
 $b=5.434\text{\AA}$
 $c=5.434\text{\AA}$
 $\alpha=90.000^\circ$
 $\beta=90.000^\circ$
 $\gamma=90.000^\circ$

space group 216
 $F\bar{4}3m$

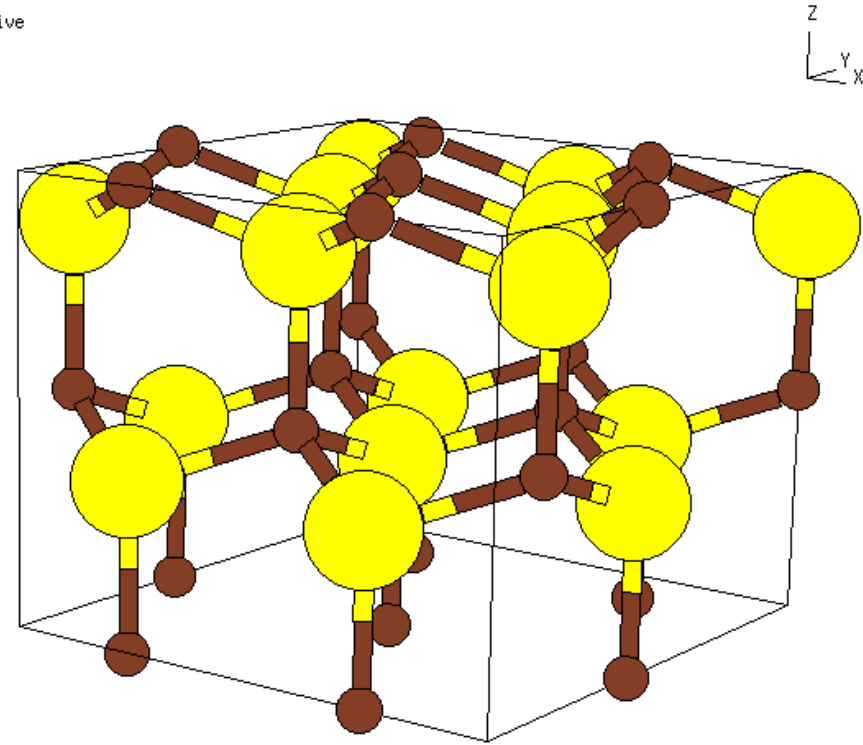


wurtzite

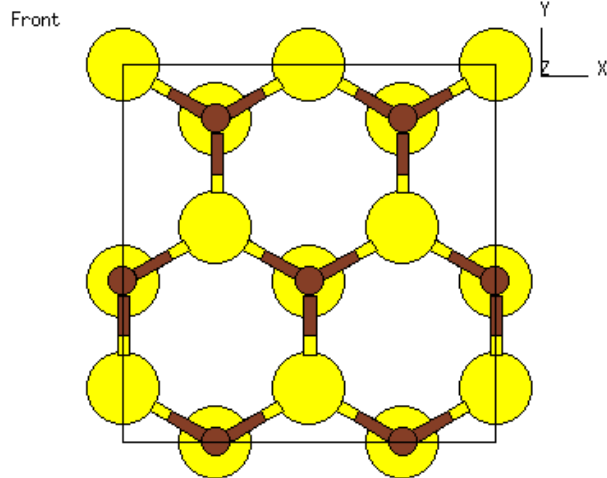
ZnS
ZnO
CdS
CdSe
GaN
AlN



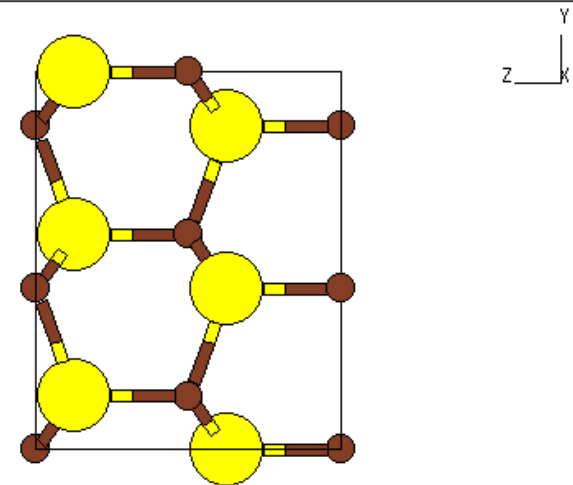
Active



Number 186



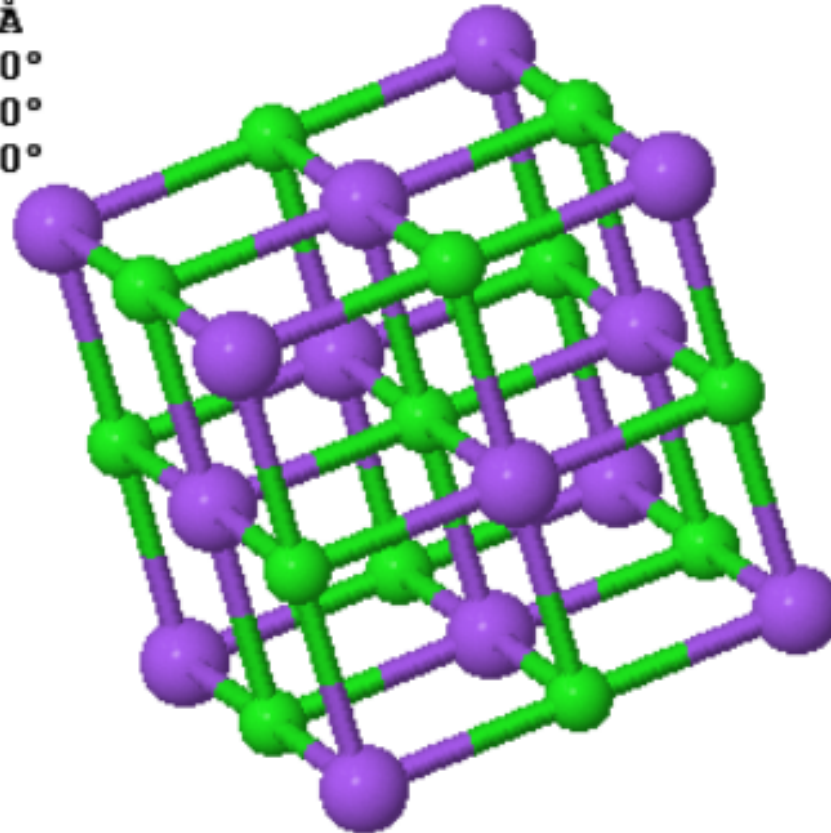
Right



There are 2 polytypes of ZnS: zincblende and wurtzite

NaCl

HM: $Fm\bar{3}m$
 $a=5.639\text{\AA}$
 $b=5.639\text{\AA}$
 $c=5.639\text{\AA}$
 $\alpha=90.000^\circ$
 $\beta=90.000^\circ$
 $\gamma=90.000^\circ$



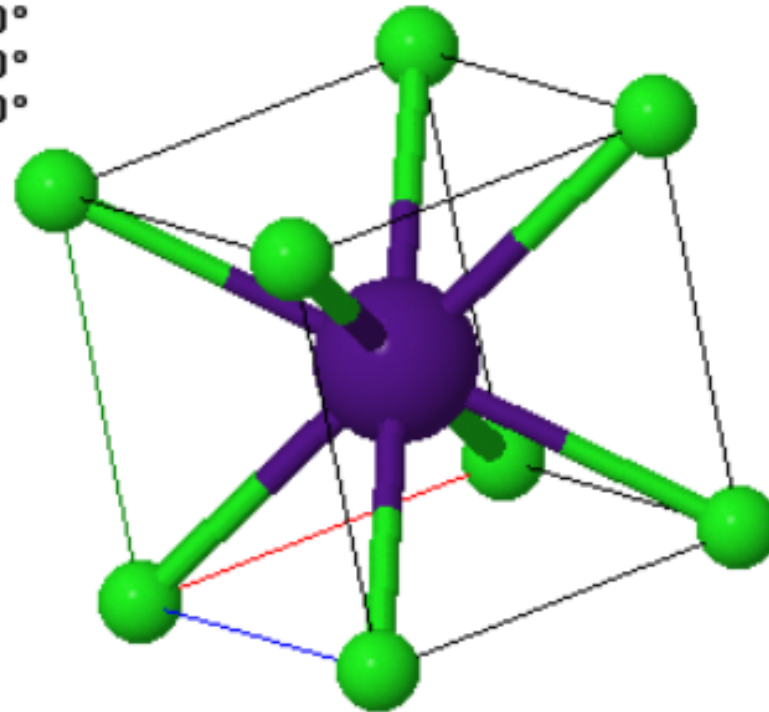
Number 225

Bravais: fcc

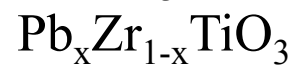
CsCl

HM: PM-3M
a=4.110Å
b=4.110Å
c=4.110Å
α=90.000°
β=90.000°
γ=90.000°

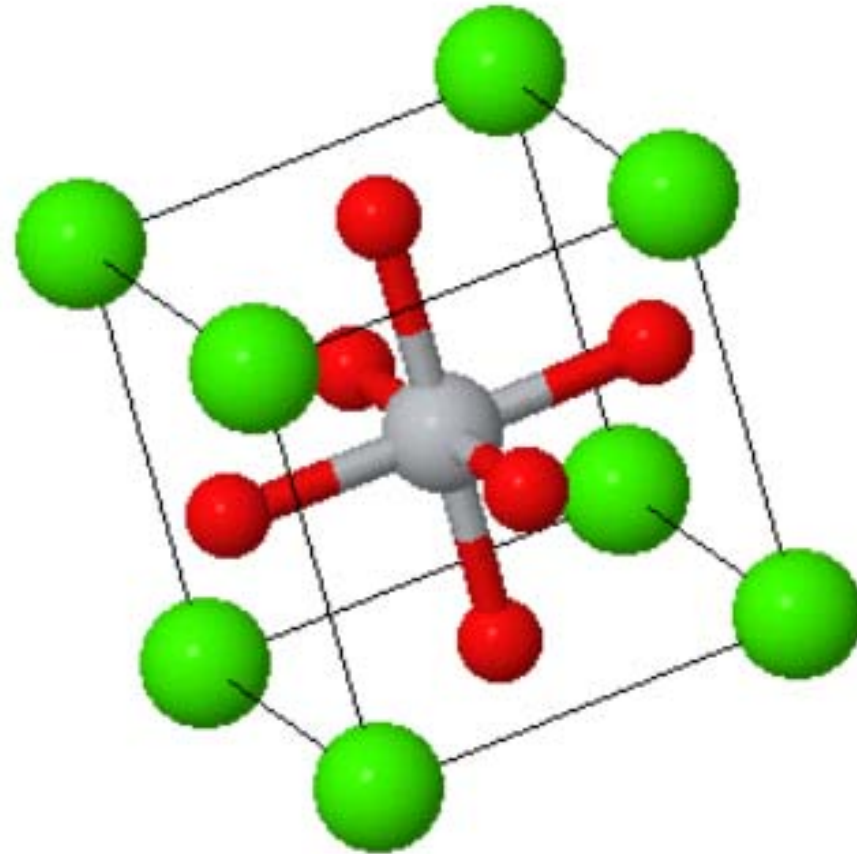
Number 221

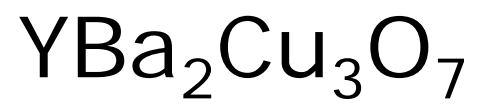


perovskite

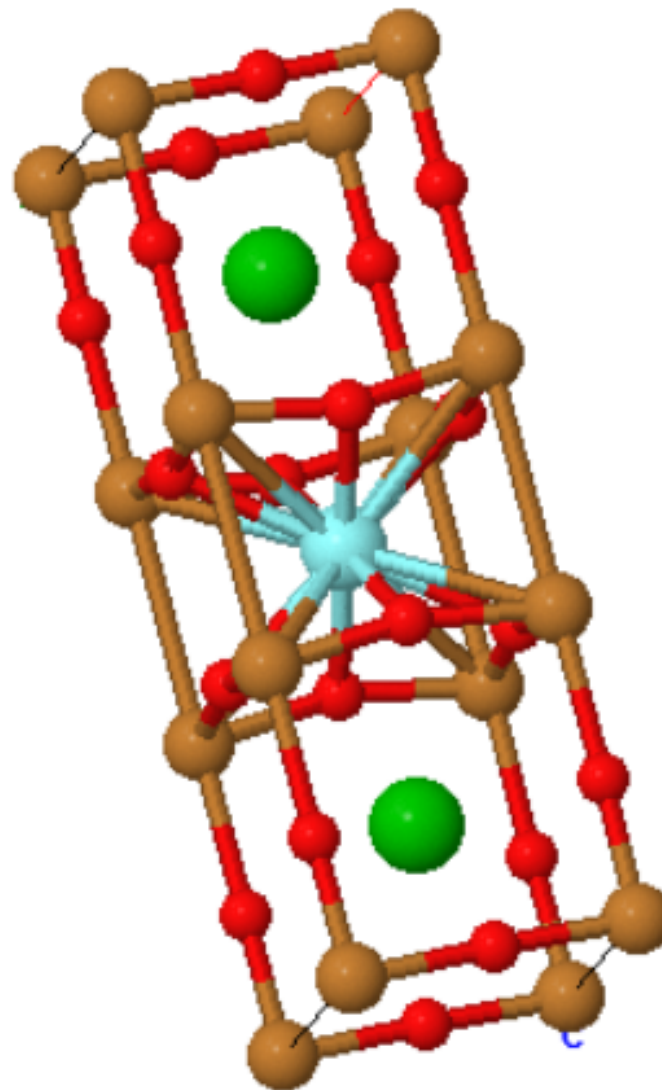


Number 221
simple cubic





HM: P m m m
a=3.820Å
b=3.885Å
c=11.683Å
 $\alpha=90.000^\circ$
 $\beta=90.000^\circ$
 $\gamma=90.000^\circ$



Number 47

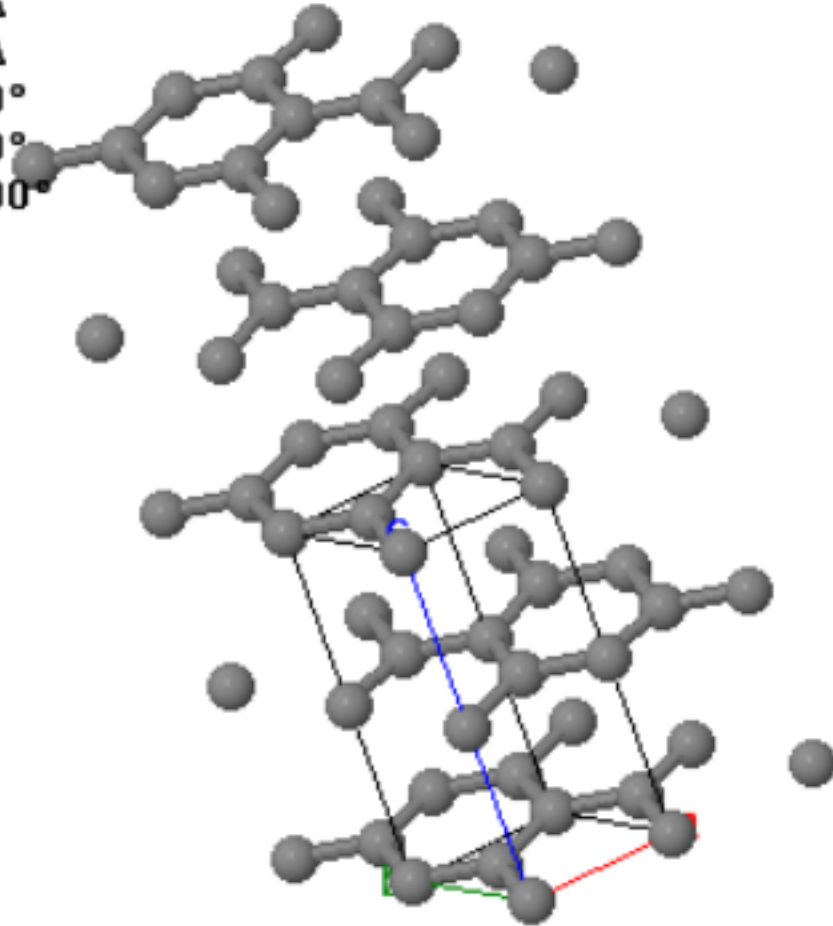
graphite

Space group 194

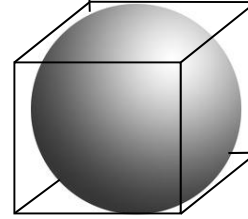
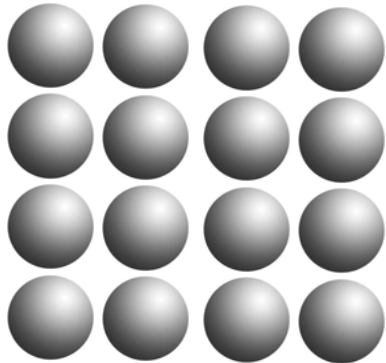
4 inequivalent C atoms in the primitive unit cell

Polytypes of carbon
graphite (hexagonal)
carbon nanotubes
diamond
rhombohedral graphite
hexagonal diamond

HM: P 63 m c
a=2.456Å
b=2.456Å
c=6.696Å
α=90.000°
β=90.000°
γ=120.000°



atomic packing density



$$\frac{\frac{4}{3} \pi (L/2)^3}{L^3} = \frac{\pi}{6} \approx 0.52$$

fcc, hcp = 0.74

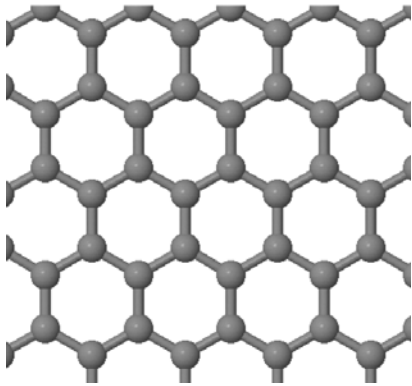
random close pack = 0.64

simple cubic = 0.52

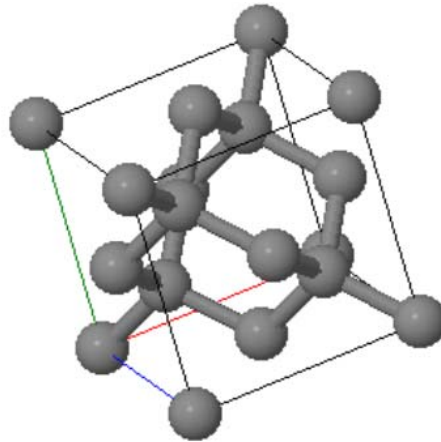
diamond = 0.34

Coordination number

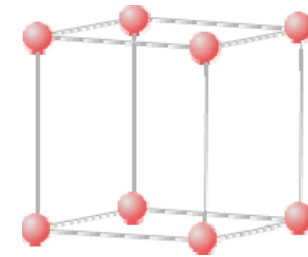
Number of nearest neighbors an atom has in a crystal



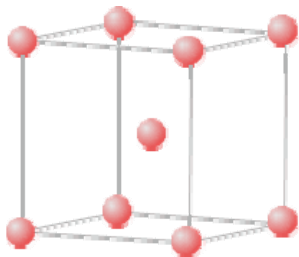
Graphene 3



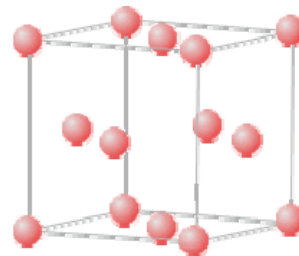
diamond 4



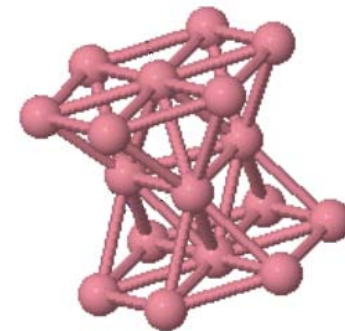
sc 6



bcc 8

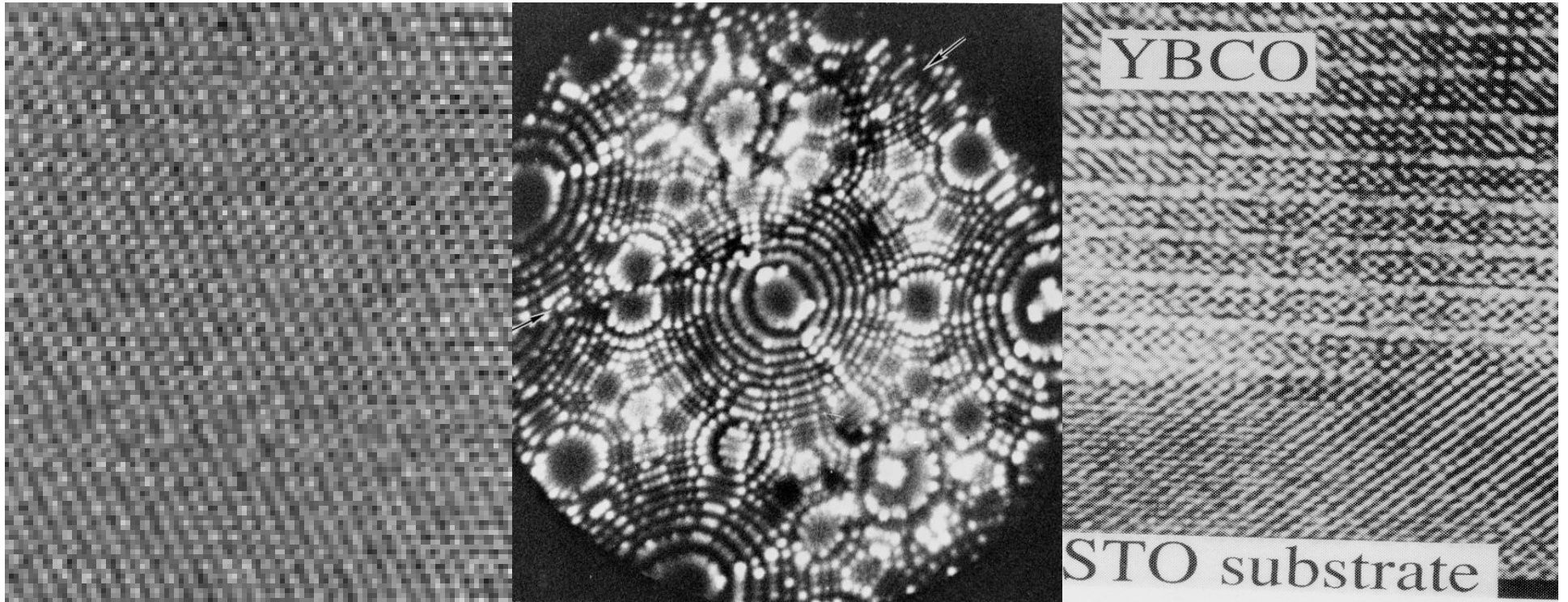


fcc 12



hcp 12

Crystal structure determination



Scanning tunneling
microscope

Field ion microscope

Transmission electron
microscope

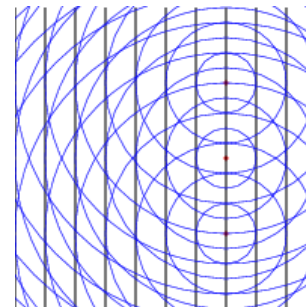
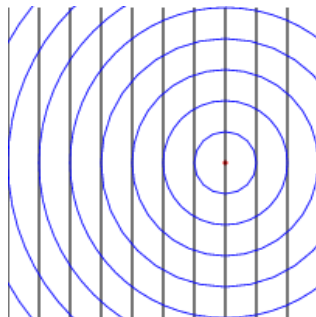
Usually x-ray diffraction is used to
determine the crystal structure

Crystal diffraction (Beugung)

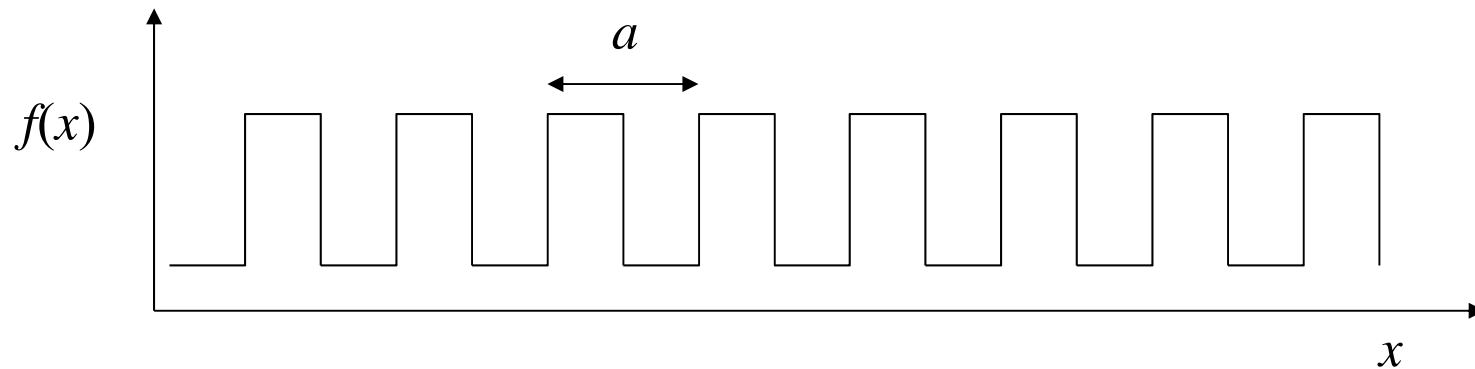
Everything moves like a wave but exchanges energy and momentum as a particle

light
sound
electron waves
neutron waves
positron waves
plasma waves

photons
phonons
electrons
neutrons
positrons
plasmons



Expanding a 1-d function in a Fourier series



Any periodic function can be represented as a Fourier series.

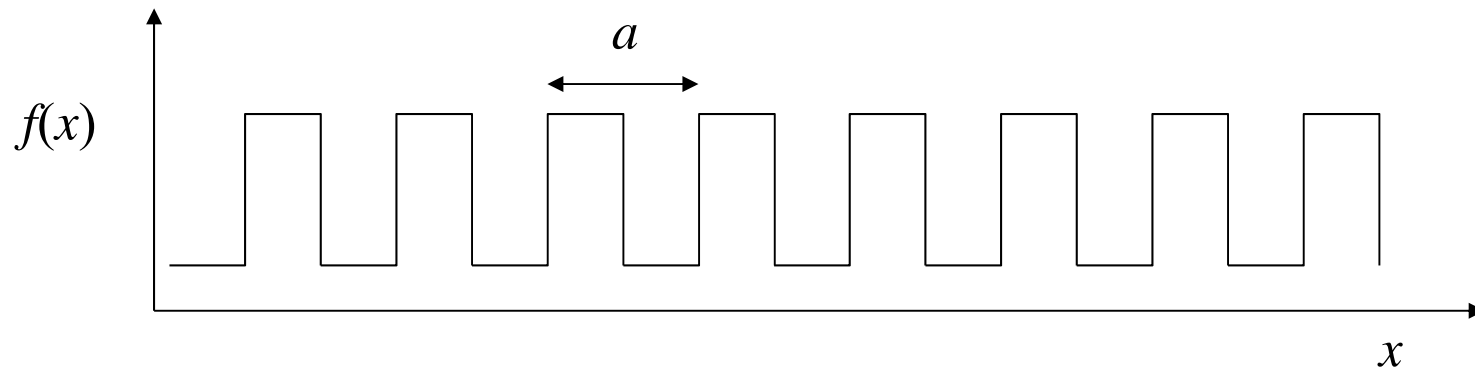
$$f(x) = f_0 + \sum_{p=1}^{\infty} c_p \cos(2\pi px/a) + s_p \sin(2\pi px/a)$$

multiply by $\cos(2\pi p'x/a)$ and integrate over a period.

$$\int_0^a f(x) \cos(2\pi px/a) dx = c_p \int_0^a \cos(2\pi px/a) \cos(2\pi px/a) dx = \frac{ac_p}{2}$$

$$c_p = \frac{2}{a} \int_0^a f(x) \cos(2\pi px/a) dx$$

Expanding a 1-d function in a Fourier series



Any periodic function can be represented as a Fourier series.

$$f(x) = f_0 + \sum_{p=1}^{\infty} c_p \cos(2\pi px/a) + s_p \sin(2\pi px/a)$$

$$\cos x = \frac{e^{ix} + e^{-ix}}{2} \quad \sin x = \frac{e^{ix} - e^{-ix}}{2i}$$

$$f(x) = \sum_{G=-\infty}^{\infty} f_G e^{iGx} \quad f_G = \frac{c_p}{2} - i \frac{s_p}{2} \quad G = \frac{2\pi p}{a}$$

For real functions: $f_G^* = f_{-G}$

reciprocal lattice vector

Fourier series in 1-D, 2-D, or 3-D

In two or three dimensions, a periodic function can be thought of as a pattern repeated on a Bravais lattice. It can be written as a Fourier series

$$f(\vec{r}) = \sum_{\vec{G}} f_{\vec{G}} e^{i\vec{G} \cdot \vec{r}}$$

Reciprocal lattice vectors
(depend on the Bravais lattice)

Structure factors
(complex numbers)

In 1-D:



$$\vec{G} = v\vec{b}$$

$$v = -\infty \dots -1, 0, 1, \dots \infty$$

$$|\vec{b}| = \frac{2\pi}{a}$$

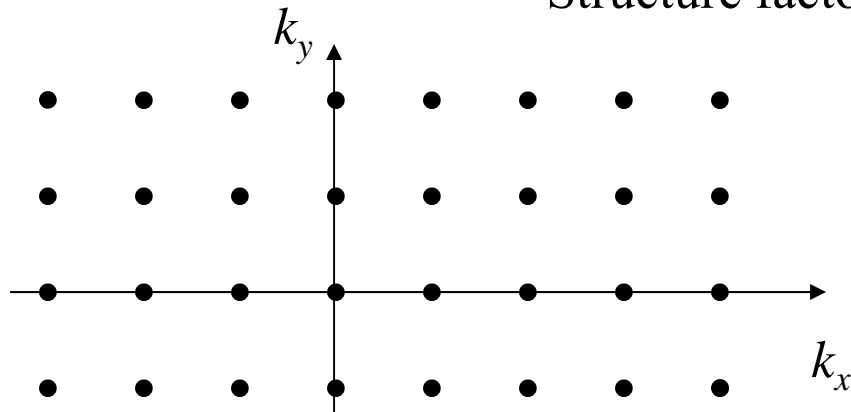
Reciprocal lattice (Reziprokes Gitter)

Any periodic function can be written as a Fourier series

$$f(\vec{r}) = \sum_{\vec{G}} f_{\vec{G}} e^{i\vec{G} \cdot \vec{r}}$$

↖ Reciprocal lattice vector G

Structure factor



$$\vec{G} = v_1 \vec{b}_1 + v_2 \vec{b}_2 + v_3 \vec{b}_3$$

v_i integers

$$\vec{b}_1 = 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)}, \quad \vec{b}_2 = 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)}, \quad \vec{b}_3 = 2\pi \frac{\vec{a}_1 \times \vec{a}_2}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)}$$

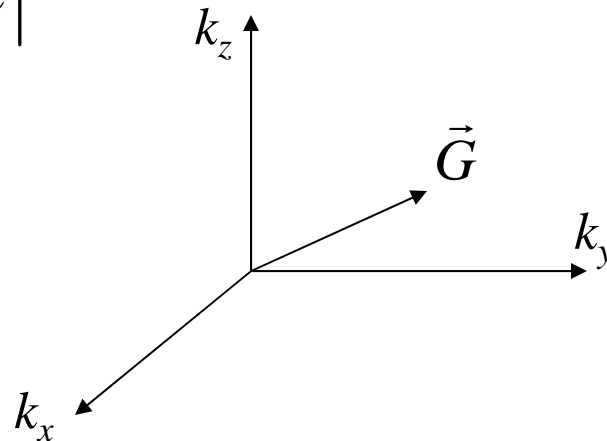
Reciprocal space (Reziproker Raum) *k*-space (*k*-Raum)

k-space is the space of all wave-vectors.

A *k*-vector points in the direction a wave is propagating.

wavelength: $\lambda = \frac{2\pi}{|\vec{k}|}$

momentum: $\vec{p} = \hbar\vec{k}$



Plane wave:

$$\exp(i\vec{G} \cdot \vec{r}) = \cos(G_x x + G_y y + G_z z) + i \sin(G_x x + G_y y + G_z z)$$